

Search for heavy lepton resonances decaying to a Z boson and a lepton in pp collisions at $\sqrt{s}=8$ TeV with the ATLAS detector

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Search for heavy lepton resonances decaying to a Z boson and a lepton in pp collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector



The ATLAS collaboration

E-mail: atlas.publications@cern.ch

ABSTRACT: A search for heavy leptons decaying to a Z boson and an electron or a muon is presented. The search is based on pp collision data taken at $\sqrt{s} = 8$ TeV by the ATLAS experiment at the CERN Large Hadron Collider, corresponding to an integrated luminosity of 20.3 fb^{-1} . Three high-transverse-momentum electrons or muons are selected, with two of them required to be consistent with originating from a Z boson decay. No significant excess above Standard Model background predictions is observed, and 95% confidence level limits on the production cross section of high-mass trilepton resonances are derived. The results are interpreted in the context of vector-like lepton and type-III seesaw models. For the vector-like lepton model, most heavy lepton mass values in the range 114–176 GeV are excluded. For the type-III seesaw model, most mass values in the range 100–468 GeV are excluded.

KEYWORDS: Hadron-Hadron Scattering

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1 Introduction

Searches for new particles often utilize decays to electrons or muons, which allow for full four-momentum reconstruction with a mass resolution better than that achievable using hadronic or semileptonic decay modes. Dilepton resonance searches led to the discovery of the J/ψ [1, 2], the Υ [3], and the Z boson [4, 5], and they have been used at the CERN Large Hadron Collider (LHC) to place strong constraints on a variety of new particles such as additional gauge bosons [6, 7]. Searches for low-mass trilepton resonances have been used to constrain lepton flavour violation in muon and τ lepton decays [8, 9].

High-mass trilepton resonances are motivated by several extensions of the Standard Model (SM). Vector-like leptons (VLL) are invoked to explain the mass hierarchy between the different lepton generations [10]. They also arise in composite Higgs models [11, 12] and models of warped extra dimensions [13, 14]. Such leptons have masses much larger than those of the SM leptons, and are defined as colourless, spin-1/2 fermions whose left- and right-handed chiral components have the same transformation properties under the weak-isospin $SU(2)$ gauge group. Another set of models predicting trilepton resonances is based on the type-III seesaw mechanism [15], which explains the origin of small neutrino masses through the introduction of heavy $SU(2)$ triplets with zero hypercharge.

This article presents a search for high-mass trilepton resonances with the ATLAS detector, using a data sample corresponding to 20.3 fb^{-1} of integrated luminosity collected

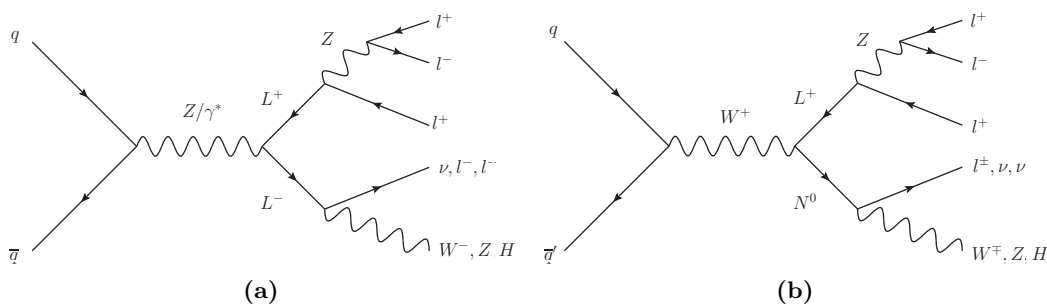


Figure 1. Feynman diagrams for the production and decay of new heavy leptons (L^\pm , N^0) to final states resulting in a trilepton resonance. Diagram (a) shows the pair production of two charged heavy leptons, and (b) shows the associated production of a charged and a neutral heavy lepton.

in pp collisions at $\sqrt{s} = 8$ TeV at the LHC. This search uses data events with at least three charged leptons (electrons or muons), two of which are consistent with originating from a Z -boson decay. Several signal regions are defined to be sensitive to the pair-production of heavy leptons that decay to SM leptons and W , Z , or H bosons. The backgrounds, dominated by SM diboson production, are estimated using Monte Carlo (MC) simulation and control regions in data, and the predictions are validated in dedicated data samples.

The results of the search are interpreted in the context of vector-like lepton [16] and type-III seesaw [17] scenarios in which the new heavy leptons decay through mixing with electrons or muons (ℓ) induced by off-diagonal Yukawa couplings. In the type-III seesaw model, the masses of the three heavy leptons are assumed to be identical. Feynman diagrams of the production and decay of the heavy leptons in both models are shown in figure 1. The heavy leptons are produced in pairs through Drell-Yan processes, with cross sections of roughly 34 fb and 844 fb for the VLL and type-III seesaw models, respectively, assuming heavy lepton masses of 200 GeV. The difference in the production cross section is due to the different gauge couplings of the models, as well as the additional neutral fermion in the type-III seesaw model. The heavy leptons decay via the mixing terms into an SM lepton and a W , Z , or H boson. The charged states L^\pm exist in both models, and have decay modes to $W\nu$, $Z\ell$, and $H\ell$; the neutral state N^0 is only present in the type-III seesaw model, and decays to $W\ell$, $Z\nu$, and $H\nu$. The charged lepton branching fractions approach $\mathcal{B}(L^\pm \rightarrow W^\pm \nu) = 50\%$, $\mathcal{B}(L^\pm \rightarrow Z\ell^\pm) = 25\%$ and $\mathcal{B}(L^\pm \rightarrow H\ell^\pm) = 25\%$ for $m_{L^\pm} \gg m_H$, in accordance with the Goldstone boson equivalence theorem [18]; at lower masses, the branching fractions to H and Z decrease as they become kinematically disfavoured. For the neutral lepton, the branching fractions to $W\ell$, $Z\nu$, and $H\nu$ are identical to those of the charged leptons to $W\nu$, $Z\ell$, and $H\ell$, respectively.

Searches for heavy leptons were previously performed at LEP, excluding vector-like leptons with masses below $m_{L^\pm} = 101.2$ GeV using the $L^\pm \rightarrow W^\pm \nu$ decay mode [19]. A search for type-III seesaw heavy leptons was performed by CMS in pp collision data at $\sqrt{s} = 7$ TeV, using non-resonant trilepton signatures to exclude seesaw fermions with masses below $m_{L^\pm} = 180$ –200 GeV, depending on the branching fractions assumed [20].

2 The ATLAS detector

The ATLAS detector [21] is a multi-purpose detector covering nearly the full solid angle¹ around the pp interaction region. The beam pipe is surrounded by the inner detector (ID), consisting of silicon pixel and microstrip detectors and a transition radiation tracker. The ID is enclosed in a superconducting solenoid providing a 2 T axial magnetic field, and performs charged particle tracking for $|\eta| < 2.5$.

The calorimeter system surrounds the solenoid, and consists of electromagnetic and hadronic components. The electromagnetic calorimeter is a lead/liquid argon (LAr) sampling calorimeter, and comprises a barrel ($|\eta| < 1.475$) and two endcaps ($1.375 < |\eta| < 3.2$). In the range $|\eta| < 2.5$, the detector is finely segmented in η to provide good spatial resolution. The hadronic calorimeter (HCAL) uses steel/scintillator tiles in the barrel ($|\eta| < 1.7$) and copper/LAr in the endcaps ($1.5 < |\eta| < 3.2$). In the forward region ($3.1 < |\eta| < 4.9$), electromagnetic and hadronic calorimetry is performed using copper/LAr and tungsten/LAr technology.

The muon spectrometer (MS) features high-precision tracking chambers interleaved with dedicated trigger chambers located in a toroidal magnetic field. The magnetic field is generated by a system of three large superconducting air-core toroid magnets, with a bending integral of about 2.5 T·m in the barrel and up to 6 T·m in the endcaps. The precision tracking is provided by monitored drift tubes ($|\eta| < 2.7$), complemented by cathode strip chambers in the forward region ($2 < |\eta| < 2.7$). Triggering is performed by resistive plate chambers in the barrel ($|\eta| < 1.05$) and thin gap chambers in the endcaps ($1.05 < |\eta| < 2.4$).

Events are recorded using a three-level trigger system. The first level, implemented in hardware, reduces the event rate to less than 75 kHz using a subset of the detector information. The second and third levels are implemented in software, and reduce the event rate to less than 400 Hz using the full detector information.

3 Object reconstruction and event selection

The data were collected during 2012 using triggers requiring either an electron or a muon with transverse momentum relative to the beam axis, p_T , greater than 24 GeV. The triggered electron or muon must also satisfy loose isolation requirements. These triggers are supplemented by triggers without isolation requirements, but with higher p_T thresholds of 60 (36) GeV for electrons (muons). Only data taken while the ID, calorimeters, and MS were functioning normally are considered. Events are required to have a reconstructed primary vertex having at least three associated tracks with $p_T > 400$ MeV, consistent with the beamspot envelope. If more than one such vertex is found, the vertex with the largest $\sum p_T^2$ of its associated tracks is chosen as the hard-scatter primary vertex.

¹ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z -axis along the beam pipe. The x -axis points from the IP to the centre of the LHC ring, and the y -axis points upward. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle θ as $\eta = -\ln \tan(\theta/2)$.

Electron candidates are selected as energy clusters within a small window of size $\Delta\eta \times \Delta\phi = 0.075 \times 0.125$ in the electromagnetic calorimeter matched to a track in the ID. They are required to fulfil tight identification criteria [22], have $|\eta_{\text{cluster}}| < 2.47$, and not be in the transition region between the barrel and the endcap calorimeter ($1.37 < |\eta_{\text{cluster}}| < 1.52$), where η_{cluster} is the pseudorapidity of the barycentre of the energy cluster. Muon candidates are selected as tracks reconstructed in the MS matched to tracks in the ID [23] and are required to satisfy $|\eta| < 2.5$. The muon momentum is determined from combining the information from the two tracks. Muons and electrons are required to have transverse momenta greater than 15 GeV and to be isolated from tracks and calorimeter energy deposits using the criteria described in ref. [24]. To ensure that the lepton track is consistent with originating from the primary event vertex, the ID track is required to satisfy $|d_0/\sigma_{d_0}| < 3$ and $|z_0 \sin \theta| < 0.5$ mm, where d_0 and z_0 are the transverse and longitudinal impact parameters of the track with respect to the primary vertex, respectively, and σ_{d_0} is the uncertainty on the transverse impact parameter. In order to ensure constant trigger efficiency as a function of lepton p_T , at least one electron or muon must have $p_T > 26$ GeV and a separation $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2}$ less than 0.2 from the triggered electron or muon. The trigger efficiency is evaluated to be larger than 95% when all offline selection criteria are applied.

Jets are reconstructed from topological clusters built from energy deposits in calorimeter cells using the anti- k_t jet algorithm [25] with a radius parameter of $R = 0.4$. The measured jet energy is calibrated using p_T - and η -dependent corrections for instrumental effects (e.g. passive material and non-compensating response of the calorimeters) derived from MC simulations and in situ techniques applied to data, and is corrected for additional pp interactions per bunch crossing (pileup) [26, 27].

After energy calibration, jets are required to have $p_T > 30$ GeV and $|\eta| < 4.5$. In order to suppress jets from pileup interactions, if a jet has $p_T < 50$ GeV and $|\eta| < 2.5$, then at least 50% of the scalar sum of the p_T of all tracks associated with the jet is required to come from tracks associated with the primary event vertex. Jets are also required to pass jet-quality selections to reject jets reconstructed from non-collision signals, such as beam-related background, cosmic rays or detector noise.

Since leptons and jet candidates can be reconstructed as multiple objects, the overlap between the various objects is resolved by applying the following procedure. If two electrons are separated by $\Delta R < 0.1$, the electron with the lower p_T is removed. If an electron and a jet are separated by $\Delta R < 0.2$, the jet is removed. If an electron and a jet satisfy $0.2 < \Delta R < 0.4$, and the jet's transverse momentum also satisfies $p_{T,\text{jet}} > 30 \text{ GeV} + 0.05 p_{T,e}$, the electron is removed. If a muon and an electron satisfy $\Delta R < 0.1$, the electron is removed. If a muon and a jet are separated by $\Delta R < 0.1$, the jet is removed if its transverse momentum satisfies $p_{T,\text{jet}} < 0.5 p_{T,\mu}$ if $p_{T,\mu} < 200$ GeV, or $p_{T,\text{jet}} < 100$ GeV if $p_{T,\mu} > 200$ GeV. Finally, if a muon and a jet not removed by the previous requirement are separated by $\Delta R < 0.3$, the muon is removed.

The missing transverse momentum, $\mathbf{p}_T^{\text{miss}}$, and its magnitude, E_T^{miss} , are calculated from the vector sum of the transverse momenta of all calibrated electrons, muons, τ leptons, jets, and all topological calorimeter clusters of energy not associated with other objects with $|\eta| < 4.9$.

Events are required to have at least three leptons (electrons or muons) passing the selection requirements above. At least one pair of leptons with the same flavour and opposite electric charge must have an invariant mass within 10 GeV of the Z boson mass, m_Z [28]. Events with four leptons consistent with the decay of two Z bosons, also within 10 GeV of m_Z , are vetoed. For the remaining events with four leptons, the lepton closest in ΔR to the Z boson candidate, referred to here as the off- Z lepton, is used to form the trilepton mass. For the range of heavy lepton masses considered in this analysis, the Z boson and the off- Z lepton tend to be collimated; hence, to improve the signal to background ratio, events where the Z candidate and the off- Z lepton are separated by $\Delta R > 3$ are vetoed.

For simulated events with an L^\pm decaying to three leptons with $p_T > 15$ GeV and $|\eta| < 2.5$, of which two originate from a Z boson and have an invariant mass within 10 GeV of m_Z , the efficiency of this event selection for the $Z + e$ ($Z + \mu$) decay channel ranges from 20% (36%) at $m_{L^\pm} = 100$ GeV to 35% (38%) at $m_{L^\pm} = 400$ GeV. The determination of the efficiency is discussed in section 8.

Since the heavy leptons are produced in pairs, in addition to the identified $L^\pm \rightarrow Z + \ell$ decay, signal events contain either a second L^\mp or an N^0 , which decays to a W , Z , or H boson and a charged or neutral lepton. A large fraction of events therefore contain a fourth lepton and/or a hadronically decaying boson. The sensitivity of the analysis is significantly improved by separating the events selected above into the following three exclusive categories:

- 4ℓ : at least four leptons are required using the same identification criteria as described above.
- $3\ell + \text{jj}$: exactly three leptons are required, along with two jets with an invariant mass satisfying $m_W - 20 \text{ GeV} < m_{\text{jj}} < 150 \text{ GeV}$, where m_W is the W boson mass [28].
- 3ℓ -only: the event does not fulfil the criteria of either the 4ℓ or the $3\ell + \text{jj}$ categories.

Subdividing the 4ℓ category based on the presence of a dijet does not significantly improve the sensitivity due to the small number of expected events with both a fourth lepton and a hadronically decaying boson. Finally, events are separated into two channels based on whether the off- Z lepton is an electron or a muon. This classification results in six independent signal regions.

The search is performed by looking for a narrowly peaked excess of events in the distributions of the mass difference defined by $\Delta m \equiv m_{3\ell} - m_{\ell^\pm \ell^\mp}$, where the invariant mass of the two leptons associated with the Z -boson decay is subtracted from the trilepton invariant mass. This reduces the impact of the lepton momentum resolution, and thus enhances the narrow resonance structure of the signal. The resulting reconstructed width in Δm is 5.9 GeV (15.5 GeV) for a mass hypothesis of $m_{L^\pm} = 120$ GeV (400 GeV), while the corresponding width in the trilepton invariant mass is 7.3 GeV (18 GeV), for final states where the off- Z lepton is an electron. For final states where the off- Z lepton is a muon, the corresponding width in Δm is 5.1 GeV (31.5 GeV) for a mass hypothesis of $m_{L^\pm} = 120$ GeV (400 GeV), while the corresponding width in the trilepton invariant mass is

6.7 GeV (33.5 GeV). The intrinsic width of the resonance is a few MeV at $m_{L^\pm} = 120$ GeV, rising to 0.5 GeV at $m_{L^\pm} = 400$ GeV.

4 Monte Carlo simulation

The analysis uses MC samples of VLL and type-III seesaw events generated with MADGRAPH 4.5.2 and 5.2.2.1 [29], respectively, using the CTEQ6L1 [30] parton distribution functions (PDF) and the AU2 underlying event tune [31]. Showering is performed with PYTHIA 8 [32]. Decays of the heavy leptons in the VLL model are performed using BRIDGE [33], while decays in the type-III seesaw samples are performed by MADGRAPH. For the type-III seesaw model, the charged and neutral heavy leptons are generated with identical masses. Vector-like lepton samples are generated for eleven mass hypotheses for $100 \text{ GeV} \leq m_{L^\pm} \leq 400 \text{ GeV}$, while the type-III seesaw samples are generated for ten mass hypotheses for $100 \text{ GeV} \leq m_{L^\pm, N^0} \leq 500 \text{ GeV}$. The cross sections for both samples are calculated at leading order (LO) in QCD.

The main backgrounds originate from SM diboson production, in particular WZ and ZZ production. Contributions from WZ (ZZ) are modelled using the SHERPA [34] MC generator version 1.4.3 (1.4.5), using the internal showering algorithm [35–37], with the CT10 [38] PDF set and normalized to the next-to-leading-order (NLO) prediction from VBFNLO-2.6.2 [39]. The generation includes up to three additional parton emissions in the matrix element. Samples of simulated events based on the NLO generator POWHEG-BOX [40] are used to derive systematic uncertainties on the shapes of distributions predicted by SHERPA. The diboson samples are showered with PYTHIA 8, and use the CT10 PDF set and AU2 underlying event tune.

Drell-Yan production in association with a photon that converts in the detector, denoted $Z + \gamma$, is modelled using SHERPA 1.4.1, also using the CT10 PDF set and including up to three additional parton emissions in the matrix element. Production of top-quark pairs in association with a W or Z boson ($t\bar{t} + V$) and triboson production ($VVV^{(*)}$) are modelled using MADGRAPH 5.1.3.33, with PYTHIA 6.426 for the parton shower and hadronization, AUET2B underlying event tune [41], and the CTEQ6L1 PDF set. The $t\bar{t} + V$ processes are normalized to the corresponding NLO cross sections [42, 43], while the $Z + \gamma$ and $VVV^{(*)}$ processes are normalized to their LO cross sections from the respective generator.

For all samples, the response of the ATLAS detector is modelled using the GEANT4 toolkit [44, 45]. Additional pp interactions in the same or nearby bunch crossings are included in the simulation by overlaying minimum-bias interactions modelled with PYTHIA 6.425 onto the hard-scatter event. The simulated events are reweighted to reproduce the distribution of the average number of pp interactions per crossing observed in data. The generator, parton shower, PDF set, underlying event tune, and accuracy of theoretical cross section for the primary MC samples used are summarized in table 1.

Process	Generator	Parton shower and hadr.	PDF set	UE tune	Cross section
VLL	MADGRAPH 4.5.2	PYTHIA 8	CTEQ6L1	AU2	LO
Seesaw	MADGRAPH 5.2.2.1	PYTHIA 8	CTEQ6L1	AU2	LO
WZ	SHERPA1.4.3	SHERPA	CT10	SHERPA	NLO
ZZ	SHERPA1.4.5	SHERPA	CT10	SHERPA	NLO
$t\bar{t} + W/Z$	MADGRAPH 5.1.3.33	PYTHIA 6.426	CTEQ6L1	AUET2B	NLO
$VVV^{(*)}$	MADGRAPH 5.1.3.33	PYTHIA 6.426	CTEQ6L1	AUET2B	LO
$Z + \gamma$	SHERPA	SHERPA	CT10	SHERPA	LO

Table 1. Summary of the primary signal and background MC samples used in this analysis. The generator, parton shower and hadronization, PDF, underlying event tune, and the order of the cross-section calculation are shown for each sample.

5 Background estimation

Standard Model processes containing three or more lepton candidates can be classified into two categories. The first category consists of events with three prompt leptons produced in the decays of electroweak gauge bosons, which are estimated using the simulated samples described above. The second consists of events where at least one reconstructed lepton arises from a misidentified jet, hadron decay, or photon conversion, and is referred to as reducible background. For muons, reducible backgrounds arise from semileptonic b - or c -hadron decays and from in-flight decays of pions or kaons. Reducible electron backgrounds can arise from semi-leptonic b - or c -hadron decays, photon conversions and misidentified hadrons or jets. Drell-Yan production of a lepton pair with an associated photon that converts in the detector and is reconstructed as an isolated lepton ($Z + \gamma$) is estimated using simulation. The remainder of the reducible background is estimated by scaling control samples in data, following a method similar to that described in ref. [46]. The control samples consist of events with one or more leptons that do not satisfy the nominal selection criteria, but instead satisfy a set of relaxed criteria, defined separately for each lepton flavour. The events are weighted with scale factors computed for each such lepton, defined as the ratio of misidentified or non-prompt lepton candidates that satisfy the nominal criteria to those which only fulfil the relaxed criteria. For electrons, the identification requirement is changed from tight to loose [22]. For muons, the requirements on the lepton isolation and on $|d_0/\sigma_{d_0}|$ are relaxed. The scale factors are measured as a function of the candidate's p_T and η in samples of data that are enriched in non-prompt and misidentified lepton candidates. The contamination from prompt leptons in the background-enriched samples is removed using simulation.

The background estimates are validated in four validation regions. The high- ΔR region consists of events where the Z boson candidate and the off- Z lepton are separated by $\Delta R > 3$. The background composition in this region is similar to that in the signal regions. The off- Z region contains events with exactly three leptons, where no opposite-sign same-flavour pair of leptons is reconstructed with an invariant mass within 20 GeV of m_Z . This region is designed to test the $Z + \gamma$ background estimate. The ZZ region consists of events with two reconstructed Z boson candidates. These three regions are completely orthogonal to the

Channel	Validation Region	Data	Background Prediction	$\frac{\text{Data}-\text{Bkgd}}{\sigma_{\text{bkgd}}}$
$Z + e$	High- ΔR	239	239 ± 14	0.0
$Z + e$	Off- Z	360	349 ± 44	+0.2
$Z + e$	ZZ	39	37 ± 2	+0.3
$Z + e$	WZ	140	133 ± 10	+0.4
$Z + \mu$	High- ΔR	302	301 ± 12	+0.1
$Z + \mu$	Off- Z	163	200 ± 8	-2.3
$Z + \mu$	ZZ	74	63 ± 3	+1.2
$Z + \mu$	WZ	222	193 ± 14	+1.5

Table 2. Summary of the number of events observed and predicted for each validation region. The uncertainty on the background prediction is the total systematic uncertainty. The difference between the observed and predicted number of events divided by the combined statistical and systematic uncertainty on the prediction is also shown.

signal selection. The fourth region is designed to test the modelling of SM WZ production. Events in this region must have exactly three leptons, zero jets, $40 < E_{\text{T}}^{\text{miss}} < 100$ GeV, and $40 < m_{\text{T}}^W < 90$ GeV, where $m_{\text{T}}^W = \sqrt{2p_{\text{T}}^{\ell} E_{\text{T}}^{\text{miss}} (1 - \cos(\Delta\phi))}$ is the transverse mass and $\Delta\phi$ is the azimuthal angle between the missing transverse momentum and the off- Z lepton with momentum p_{T}^{ℓ} . This region is not completely orthogonal to the signal regions, but signal processes are expected to account for less than 3% of the expected event yield for type-III seesaw leptons with $m_{L\pm} > 160$ GeV.

The expected and observed numbers of events are given in table 2 for all validation regions, separately for the $Z + e$ and the $Z + \mu$ channels. The largest difference is seen in the off- Z region in the $Z + \mu$ flavour channel, where there is a deficit in the data corresponding to 2.3 standard deviations (σ). The region is dominated by contributions from ZZ , where only three leptons pass the selection requirements and no same-flavour, opposite-sign lepton pair is reconstructed with invariant mass within 20 GeV of m_Z . In the other seven regions, agreement better than 1.5σ is observed. Figure 2 shows the Δm distributions for the high- ΔR and ZZ validation regions in the $Z + e$ and the $Z + \mu$ flavour channels.

6 Systematic uncertainties

Systematic uncertainties are assigned to the signal and background predictions derived from simulation to account for possible modelling inaccuracies. Sources of uncertainty affecting all simulated signal and background processes are the cross sections of SM processes, trigger efficiencies, lepton energy scales and resolutions (LES/LER), jet energy scale and resolution (JES/JER), lepton reconstruction and selection efficiencies, MC statistical uncertainties, and luminosity. The cross-section uncertainties, evaluated for the SM background samples, include renormalization and factorization scale and PDF uncertainties. The scale uncertainties are determined by varying the renormalization and factorization scales up and down by factors of two. PDF uncertainties are obtained using the PDF4LHC working group recommendations [47]. Scale and PDF uncertainties are added in quadra-

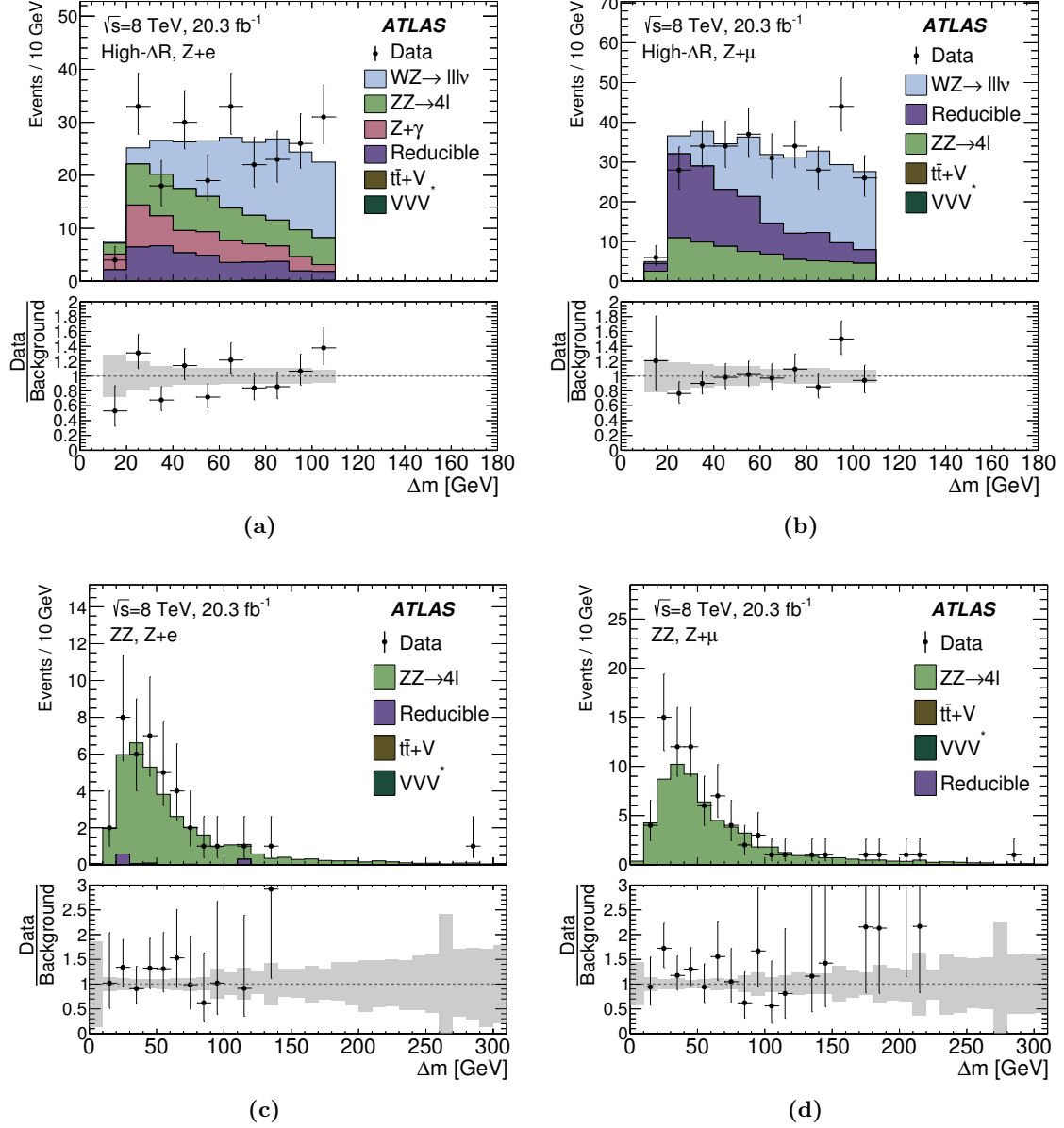


Figure 2. The $\Delta m = m_{3\ell} - m_{\ell+\ell^-}$ distributions for the high- ΔR validation region (a) and (b) and the ZZ validation region (c) and (d). The figures (a) and (c) show the distributions for the $Z + e$ final states, while (b) and (d) show the $Z + \mu$ final states. The error bars on the data points represent statistical uncertainties, and the shaded band represents the systematic uncertainty on the background prediction.

ture to obtain the total uncertainty on the inclusive cross section. For the dominant WZ and ZZ backgrounds, the resulting theoretical uncertainty on the NLO predictions from VBFNLO are 7.6% and 4.3%, respectively. A further uncertainty is assigned to the WZ and ZZ backgrounds to account for potential generator-level mismodelling of the shape of the Δm spectrum. The uncertainty is the difference between the predictions from SHERPA and POWHEG-BOX, symmetrized around the value from SHERPA.

For the $Z+\gamma$ backgrounds, an uncertainty of 30% is assigned to the modelling of prompt photons converting in the inner detector, based on comparisons of conversion processes in $Z \rightarrow ee$ events between data and simulation. The reducible backgrounds are assigned an uncertainty related to the data-driven scaling procedure described in section 5, primarily due to the extrapolation of the scale factors from the measurement sample to the signal regions and to the correction for the presence of prompt leptons in the background-enriched samples. The uncertainties on the electron factors range from 24% to 30% as a function of p_T , and the uncertainties on the muon factors range from 25% to 50%.

The uncertainties on the lepton reconstruction and selection efficiencies, energy scales, and energy resolutions [23, 48, 49] affect all simulated backgrounds, with combined uncertainties of 1% to 2% on the normalizations. The jet energy scale and resolution uncertainties [27, 50] only significantly affect the $3\ell+jj$ signal regions, with a total uncertainty of 3%. Statistical uncertainties due to the finite number of events in the MC samples range from 1% to 5%. The luminosity uncertainty is 2.8%, and is derived using the same methodology as that described in ref. [51]. In total, the systematic uncertainties on the background predictions in each signal region range from 6% to 9%.

The largest sources of uncertainty affecting the signal predictions are the lepton reconstruction and selection efficiencies, the luminosity, and, for the $3\ell+jj$ category, the jet energy scale and resolution. The total uncertainties on the signal normalizations range from 3% to 7% depending on the signal region and m_{L^\pm} .

7 Signal and background model

The numbers of signal and background candidate events in data are determined from an unbinned maximum-likelihood fit of a combination of signal and background models to the Δm distributions in each signal region. The details of the fit procedure and the models are described below.

The signal and background processes are modelled with probability density functions (p.d.f.s). The parameters of the p.d.f.s are determined from fits to the background estimates described in section 5. The fit to data using the combined signal and background model is performed simultaneously on the three categories for each of the two flavour channels. In each signal region, the normalization of the dominant background (ZZ or WZ) is a free parameter in the fit. The normalizations of all other backgrounds are constrained to fluctuate according to Gaussian probability distributions with mean and width values equal to the estimates and the total uncertainties before fitting. The uncertainties on the p.d.f. parameters are incorporated as Gaussian-distributed nuisance parameters.

The VLL and type-III seesaw signal models are parameterized separately as the sum of a Voigtian function (the convolution of a Breit-Wigner and a Gaussian function) for the trilepton resonance peak and a Landau distribution for the combinatorial part of the signal, where the three reconstructed leptons do not originate from the same L^\pm decay. The signal parameterization at a certain heavy lepton mass m_{L^\pm} is given as a function of Δm by the following expression:

$$\mathcal{S}(m_{L^\pm}) = f_V F_V(\Delta m; \Gamma_V, m_V, \sigma_V) + (1 - f_V) F_L(\Delta m; m_L, \sigma_L), \quad (7.1)$$

where f_V denotes the fraction of events in the resonance peak (Voigtian function), Γ_V , m_V and σ_V are the width, mean, and Gaussian smearing terms of the Voigtian function F_V , and m_L and σ_L are the parameters of the Landau distribution F_L . The six parameters are determined at each simulated mass point by fitting $\mathcal{S}(m_{L^\pm})$ to the simulated Δm distributions, separately for the two flavours. For mass points m_{L^\pm} that lie between those assumed in the MC samples, the parameters of the signal templates are obtained by linearly interpolating the fitted values determined at the nearest simulated mass points. The fraction of events in the Voigtian part of the signal, f_V , is $\sim 60\%$ ($\sim 70\%$) at 120 GeV, decreasing to $\sim 58\%$ ($\sim 55\%$) at 400 GeV for the type-III seesaw signal (VLL). The uncertainties on the fit parameters of the signal p.d.f. are incorporated as Gaussian-distributed nuisance parameters in the combined fits to data.

The combined background model consists of five different p.d.f.s, corresponding to WZ , ZZ , $Z + \gamma$, reducible, and the sum of the $t\bar{t} + V$ and triboson backgrounds. The leading WZ and ZZ backgrounds are both modelled with a modified Bukin function [52], a three-parameter function designed to model peaks with asymmetric tails using the convolution of a Gaussian and an exponential function. To mitigate the impact of MC statistical uncertainties, the parameterizations for the 4ℓ and 3ℓ -only categories are determined from the combination of all three categories; for the $3\ell + jj$ category, a separate parameterization is used to account for possible kinematic effects from the two additional jets. The uncertainty on the shape of the Δm distribution predicted by the generator is taken into account with a Gaussian-distributed nuisance parameter multiplying a template given by the difference between the p.d.f.s determined from SHERPA and POWHEG-BOX. Finally, in the $3\ell + jj$ and 3ℓ -only categories, the ratio of normalizations of the WZ and ZZ backgrounds is fixed to the prediction from MC simulation, due to the inability of the fit to resolve the similar shapes of the Δm distributions. In the 4ℓ category, the contribution from WZ events is negligible, so only the ZZ background normalization is left as a free parameter.

The most important remaining backgrounds are due to reducible processes and $Z + \gamma$ production. The reducible background is parameterized with a Landau distribution, determined from the data-driven estimate described above. Due to the low number of events in the reducible background estimates, the corresponding shape is obtained by fitting the combined distribution of all six signal regions.

The $Z(\ell\ell) + \gamma$ background contributes significantly only to the $Z + e$, 3ℓ -only category. This background is modelled with the sum of a Landau distribution and a Gaussian function. The normalizations of the reducible and $Z + \gamma$ backgrounds are constrained to the expected values, each with a Gaussian-distributed uncertainty of 30%.

	$Z + e$			$Z + \mu$		
Process	4l SR	3l + jj SR	3l-only SR	4l SR	3l + jj SR	3l-only SR
	Before combined background-only fit					
ZZ	10.9 \pm 0.6	11.7 \pm 0.8	91 \pm 5	21.4 \pm 1.1	7.5 \pm 0.6	90 \pm 5
WZ	0.08 \pm 0.01	35.3 \pm 3.1	337 \pm 28	—	46 \pm 4	480 \pm 40
$Z + \gamma$	—	2.3 \pm 0.8	35 \pm 11	—	—	—
Reducible	—	1.6 \pm 0.5	38 \pm 14	1.5 \pm 0.3	8.8 \pm 3.0	79 \pm 22
$t\bar{t} + V, VVV^{(*)}$	1.2 \pm 0.2	7.8 \pm 1.7	2.3 \pm 0.4	1.5 \pm 0.2	9.5 \pm 2.1	3.3 \pm 0.5
Total Background	12.2 \pm 0.7	59 \pm 4	504 \pm 34	24.4 \pm 1.2	72 \pm 6	650 \pm 50
	After combined background-only fit					
ZZ	15 \pm 4	13.4 \pm 2.3	107 \pm 9	22 \pm 5	10.1 \pm 1.6	88 \pm 8
WZ	0.08 \pm 0.03	39 \pm 6	393 \pm 28	0.02 \pm 0.02	56 \pm 9	460 \pm 40
$Z + \gamma$	—	2.2 \pm 0.8	34 \pm 11	—	—	—
Reducible	—	1.8 \pm 1.2	37 \pm 13	1.8 \pm 0.9	10.2 \pm 2.8	92 \pm 24
$t\bar{t} + V, VVV^{(*)}$	1.1 \pm 0.3	7.5 \pm 1.7	2.5 \pm 0.6	1.5 \pm 0.4	9.1 \pm 2.1	3.3 \pm 0.8
Total Background	16 \pm 4	64 \pm 7	574 \pm 34	25 \pm 5	85 \pm 10	640 \pm 40
Data	16	64	573	25	86	651

Table 3. Observed and expected number of events in the six signal regions, before and after the combined unbinned maximum-likelihood fit. The pre-fit uncertainties represent the total systematic uncertainties on the background estimates. The post-fit uncertainties are determined by the maximum-likelihood fit.

Finally, the background contributions from triboson and $t\bar{t} + W/Z$ production are modelled together using a Landau distribution, due to their similar shape and small contribution to the total background estimate. The normalization in the combined fit is also constrained to the expected value from simulation, with a 30% Gaussian-distributed uncertainty.

8 Results

The total number of events observed in each signal region is shown in table 3, along with the estimated backgrounds before and after fitting the total background model to the data. The corresponding Δm distributions for the pre-fit background estimates and the data are shown in figure 3. The signals expected for the VLL model with $m_{L^\pm} = 140$ GeV and the type-III seesaw model with $m_{L^\pm} = 300$ GeV are superimposed on the background as illustrative examples. The data agree with the background expectation in all cases, and no clear peak indicating resonant trilepton production is seen in any of the signal regions.

Good agreement is seen between the pre-fit and post-fit normalizations for the 4 ℓ and 3 ℓ +jj categories in the $Z + \mu$ flavour channel. The largest change in normalization due to the fit is in the 4 ℓ category for the $Z + e$ flavour channel, where the fitted ZZ normalization exceeds the prediction by 35%. The WZ and ZZ normalizations increase by roughly 15% in the 3 ℓ +jj and 3 ℓ -only categories in the $Z + e$ flavour channel, and 30% in the 3 ℓ +jj

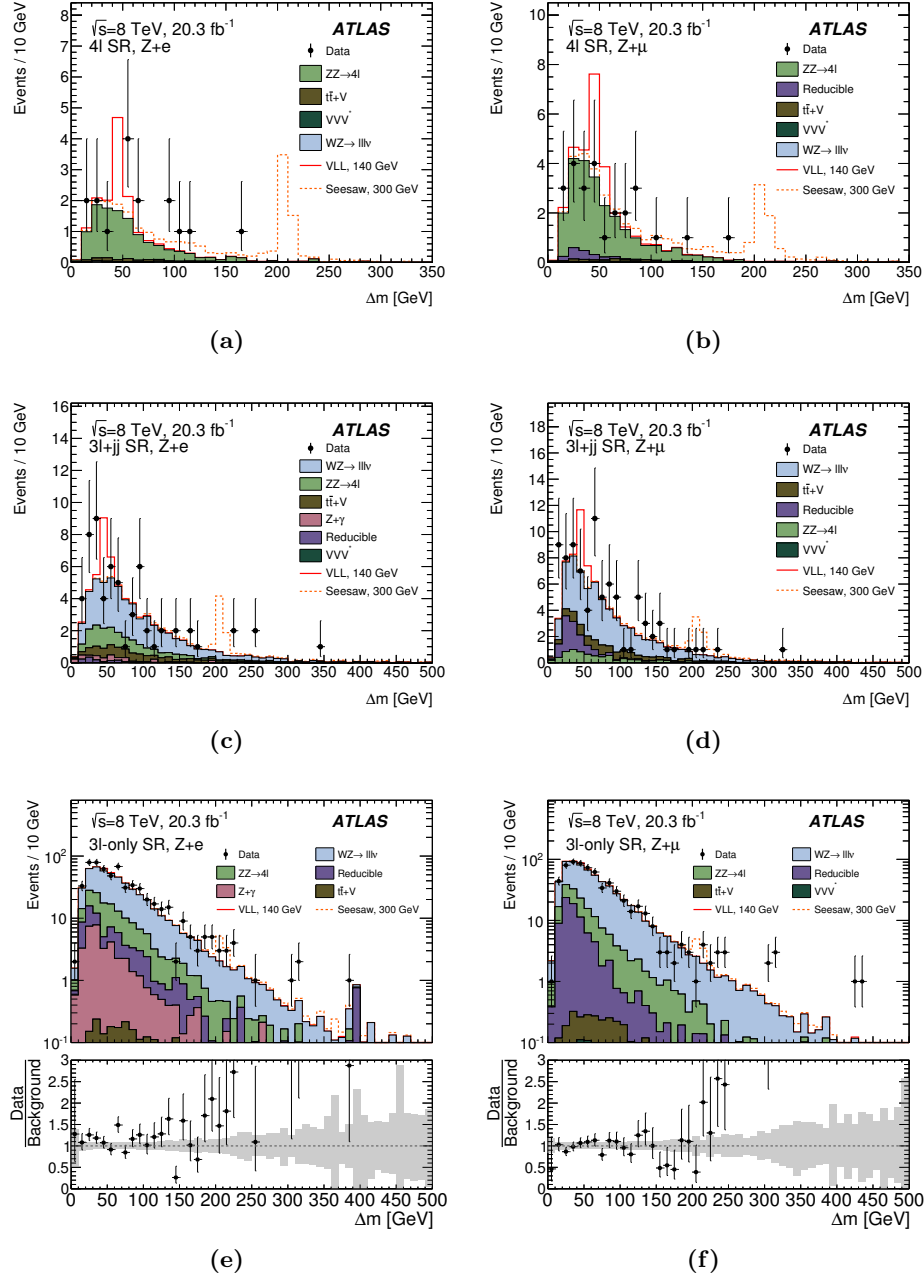


Figure 3. The $\Delta m = m_{3\ell} - m_{\ell+\ell^-}$ distributions for the 4 ℓ (top), 3 ℓ +jj (middle), and 3 ℓ -only (bottom) categories, divided into the $Z + e$ (left) and $Z + \mu$ (right) flavour channels. The observed data are shown as black points, while the pre-fit background expectations are shown in the coloured histograms. Also shown are examples for signal contributions for a 140 GeV L^\pm in the VLL model and a 300 GeV L^\pm in the type-III seesaw model. The error bars on the data points represent statistical uncertainties, and the shaded bands represent the systematic uncertainties on the background predictions.

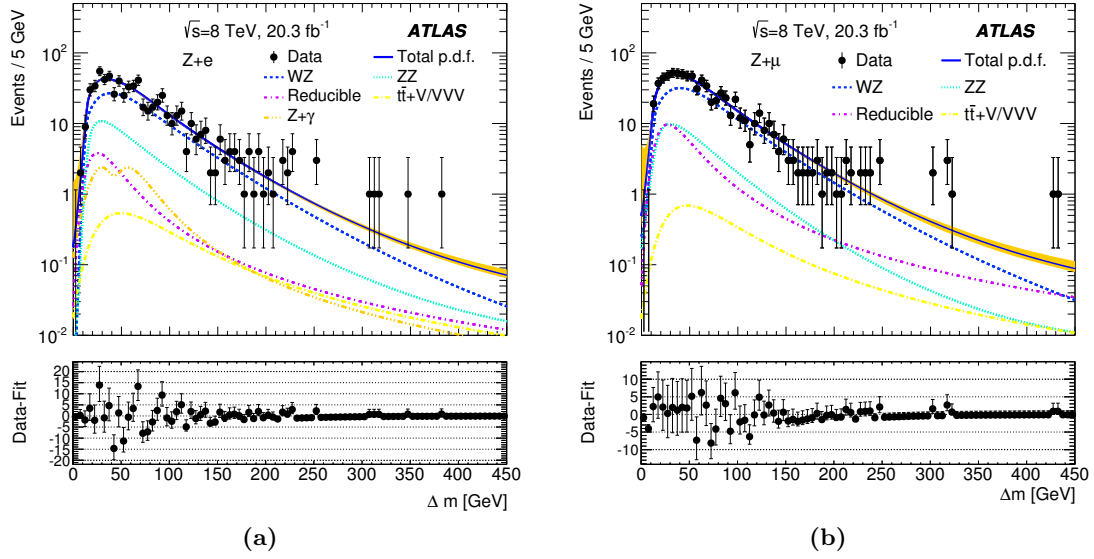


Figure 4. Projections onto the Δm variable of the background-only unbinned maximum-likelihood fits, shown superimposed on the data with the three categories in each flavour channel added together. The $Z + e$ flavour channel is shown in (a), and the $Z + \mu$ channel is shown in (b). The contributions of the separate background components to the total background-only fit are also shown. The error bars on the data points represent statistical uncertainties. Good agreement is observed between the background model and the data.

category in the $Z + \mu$ flavour channel. The projections of the fit results in the background-only hypothesis are shown in figure 4 for the combination of the three categories in each flavour channel.

The data are well described by the combined fit to the three categories in each flavour channel. The consistency of the data with the background-only hypothesis is evaluated by scanning the local p_0 -value for the Δm distribution in 3 GeV intervals for signal mass hypotheses in the range 100 – 400 GeV for the VLL model, and 100 – 500 GeV for the seesaw model, using the unbinned maximum-likelihood fit described in section 7 with the signal strength set to zero. The p_0 -value, which corresponds to the probability to observe at least as many events as observed in the present measurement assuming the background-only hypothesis, is calculated using the frequentist hypothesis test based on the profile likelihood ratio test statistic and approximated with asymptotic formulae [53]. The minimum p_0 -value is $p_0 = 0.02$ at a mass of 183 GeV for the $Z + e$ flavour channel, and $p_0 = 0.05$ at a mass of 109 GeV for the $Z + \mu$ flavour channel.

Since no significant excess above the background expectation is observed, the fit model is used to derive 95% confidence level (CL) exclusion limits on the heavy lepton pair-production cross section, σ , using the CL_s method [54]. The limits are shown for the VLL model in figure 5, and for the type-III seesaw model in figure 6, evaluated in the same 3 GeV intervals as the p_0 -values. The VLL model is excluded for electron-only mixing in the heavy lepton mass ranges 129–144 GeV and 163–176 GeV, with an expected exclusion in the

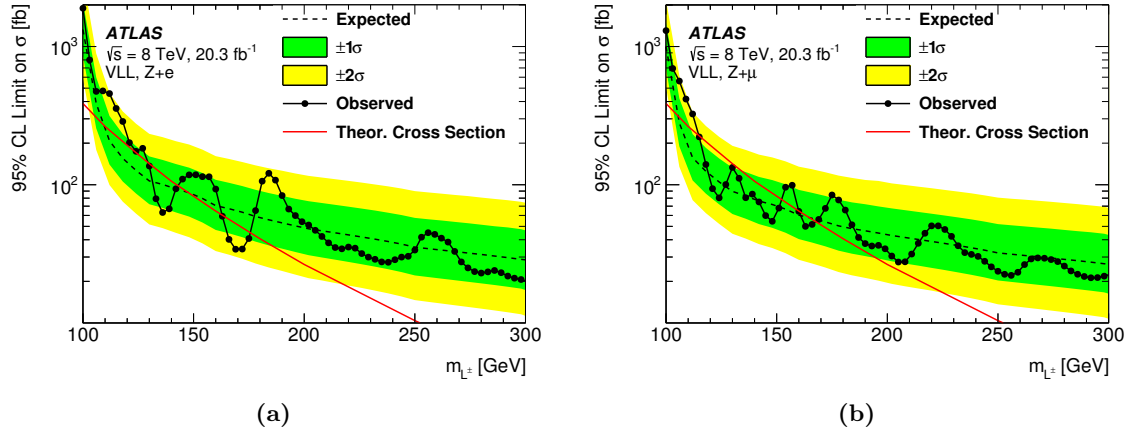


Figure 5. 95% CL upper limits on the vector-like lepton cross section. The left (right) plot shows the limits assuming 100% branching fraction to e/ν_e (μ/ν_μ). The solid line shows the observed limit. The dashed line shows the median expected limit for a background-only hypothesis, with green and yellow bands indicating the expected fluctuations at the $\pm 1\sigma$ and $\pm 2\sigma$ levels. The limit is evaluated in 3 GeV intervals.

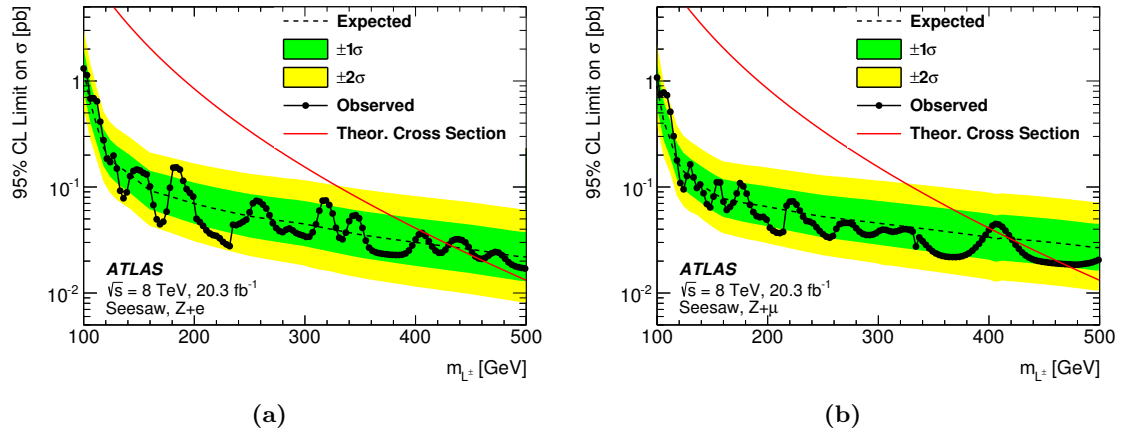


Figure 6. 95% CL upper limits on the type-III seesaw production cross section. The left (right) plot shows the limits assuming 100% branching fraction to e/ν_e (μ/ν_μ). The solid line shows the observed limit. The dashed line shows the median expected limit for a background-only hypothesis dataset, with green and yellow bands indicating the expected fluctuations at the $\pm 1\sigma$ and $\pm 2\sigma$ levels. The limit is evaluated in 3 GeV intervals.

range 109–152 GeV. The corresponding observed (expected) exclusion for the muon-only mixing scenario is 114–153 GeV and 160–168 GeV (105–167 GeV). The significantly higher production cross sections for the type-III seesaw model lead to an observed (expected) exclusion in the electron-only mixing scenario in the heavy lepton mass range 100–430 GeV (100–436 GeV). For the muon-only mixing scenario, the observed exclusion is in the ranges 100–401 GeV and 419–468 GeV, while the expected exclusion is 100–419 GeV.

The constraints shown in figures 5 and 6 are relevant to the specific VLL and type-III seesaw models considered, and are not necessarily applicable to other scenarios predicting trilepton resonances with an intermediate Z boson. A more model-independent observable is the *visible cross section*, σ_{vis} , defined as the number of observed events with $Z+\ell$ -induced trilepton resonances for a given resonance mass divided by the integrated luminosity of the data sample, 20.3 fb^{-1} . The 95% CL upper limits on σ_{vis} , denoted σ_{95}^{vis} , are derived from a fit to each flavour channel with $f_V = 1$, i.e. using only the peak component of the signal. The results for the two flavour channels, derived from the inclusive event selection without dividing the events into the three categories, are shown in figure 7.

The limits on σ_{vis} can be used to test specific models after taking into account the model's acceptance with respect to a fiducial volume, \mathcal{A} , and reconstruction and selection efficiency of events within the fiducial volume, ϵ_{fid} . The 95% CL upper limit on the cross section for the model is given by:

$$\sigma_{95} = \frac{\sigma_{95}^{\text{vis}}}{\mathcal{A} \times \epsilon_{\text{fid}}}. \quad (8.1)$$

The acceptance \mathcal{A} is defined as the probability for generated signal events to lie within a fiducial volume defined by the kinematics of the generated leptons. The leptons are considered at *particle level*, i.e. after parton shower and hadronization and with lifetimes longer than 10^{-11} s , and are *dressed*, including the contributions from radiated photons within a cone of $\Delta R = 0.1$. The fiducial volume requires that events contain an L^\pm decaying to a prompt electron or muon and a Z boson that then decays to electrons or muons. The three leptons from the L^\pm decay are required to have $p_T > 15 \text{ GeV}$ and lie within $|\eta| < 2.5$, with at least one lepton satisfying $p_T > 26 \text{ GeV}$. Two of the leptons must form a same-flavour opposite-sign pair with a mass within 10 GeV of m_Z , and the Z boson and the off- Z lepton must be separated by $\Delta R < 3$. The events are divided into flavour channels according to the flavour of the off- Z lepton. For the VLL and type-III seesaw models used in this analysis, the acceptance of events containing an $L^\pm \rightarrow Z(\ell\ell)\ell$ decay to fall within the fiducial volume is in the range 60%–65% for most of the mass range, decreasing at higher masses due to the cut on the ΔR between the Z boson and the off- Z lepton. The acceptance decreases at low masses due to the lepton p_T requirement, reaching 30%–35% at $m_{L^\pm} = 100 \text{ GeV}$.

For type-III seesaw and VLL events within the fiducial volume, ϵ_{fid} ranges from 20% to 49% if the other heavy lepton decays to a neutrino and a W , Z , or H boson.² If the other heavy lepton decays to an electron or a muon, the efficiency is 10%–20% lower, due to the increased probability of incorrectly selecting the off- Z lepton. The event selection efficiencies for the type-III seesaw model in scenarios where the second heavy lepton decays to a W boson are shown in figure 8 as a function of m_{L^\pm} ; the efficiencies for scenarios where the second heavy lepton decays to a Z or H boson and for the VLL model are consistent with these efficiencies within the statistical uncertainties.

²Note that the quoted efficiencies are dependent on the modelling of the polarization of the Z bosons, due to the requirements imposed on lepton isolation and separation.

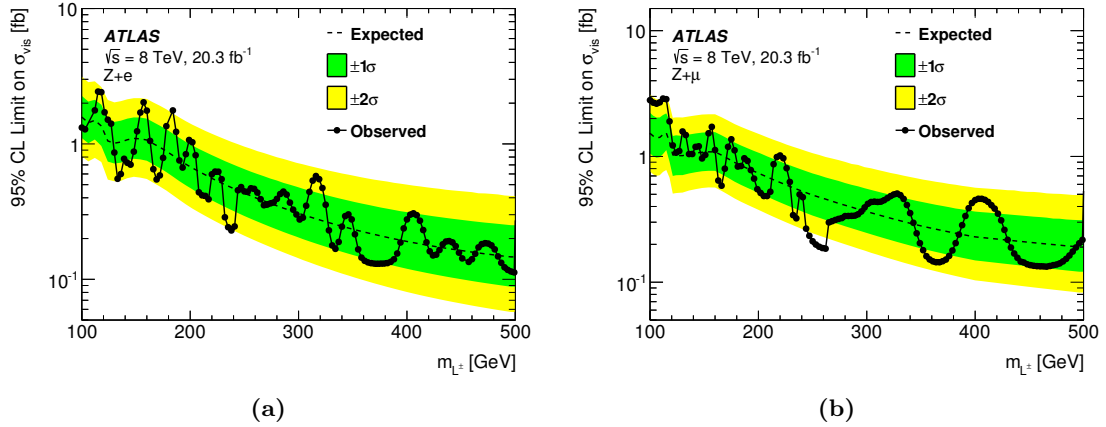


Figure 7. Upper limits at 95% CL on σ_{vis} for the $Z + e$ (left) and $Z + \mu$ (right) flavour channels, derived without dividing events into the three categories. The limits are evaluated in 3 GeV intervals.

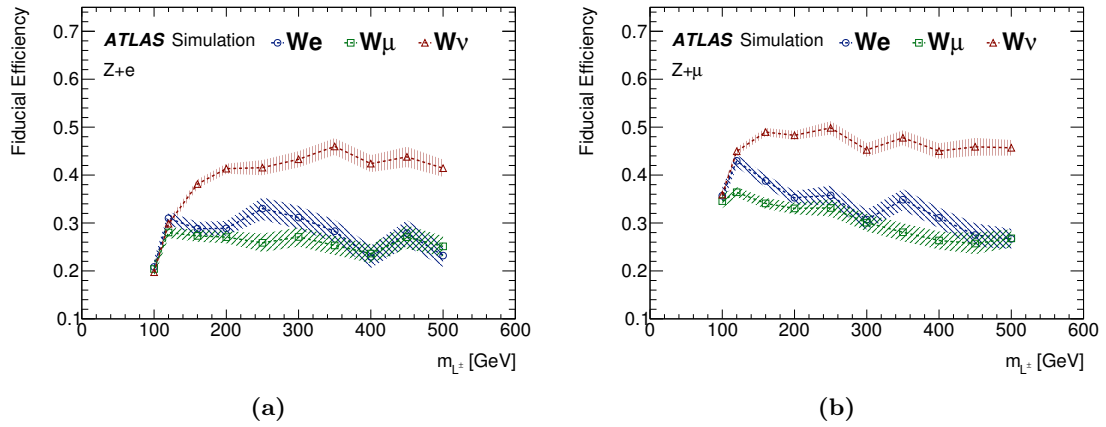


Figure 8. Efficiencies for reconstructing and correctly identifying the $L^\pm \rightarrow Z(\ell\ell)\ell^\pm$ decay in events within the fiducial volume for the type-III seesaw model. The left (right) plot shows the efficiencies for events containing a $L^\pm \rightarrow Z(\ell\ell)e$ ($L^\pm \rightarrow Z(\ell\ell)\mu$) decay. The decay of the second heavy lepton is specified in the legend. The shaded bands show the statistical uncertainty.

9 Conclusion

A search for trilepton resonances decaying to a Z boson and an electron or a muon has been presented. The search is based on pp collision data taken at $\sqrt{s} = 8$ TeV corresponding to an integrated luminosity of 20.3 fb^{-1} , collected by the ATLAS experiment at the CERN Large Hadron Collider. Events are selected requiring at least three electrons or muons with high transverse momentum, with two of the leptons consistent with the decay of a Z boson. The events are categorized based on the presence or absence of additional leptons or dijet pairs in the event consistent with the decay products of a second heavy lepton, and separated into channels based on the flavour of the lepton associated with the Z boson to form a heavy lepton decay candidate. Using the difference between the trilepton and the Z boson

candidate masses, a search for a narrow resonance is performed in each of these categories using a maximum-likelihood fit of parameterized signal and background shapes to the data. No significant excess above Standard Model predictions is observed, and 95% CL limits on the production of trilepton resonances beyond the Standard Model are derived. The results are interpreted in the context of two models of new heavy leptons decaying to three charged leptons. In the vector-like lepton model, new heavy charged leptons are excluded in the mass range 129–176 GeV (114–168 GeV) for electron-only (muon-only) mixing, except for the interval 144–163 GeV (153–160 GeV). In the type-III seesaw model, the corresponding exclusion is in the mass range 100–430 GeV (100–468 GeV) for electron-only (muon-only) mixing, except for the interval 401–419 GeV in the muon case. Limits are also set on the visible cross section of trilepton resonance productions, and fiducial efficiencies are derived to facilitate model testing.

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G. Aad⁸⁵, B. Abbott¹¹³, J. Abdallah¹⁵¹, O. Abdinov¹¹, R. Aben¹⁰⁷, M. Abolins⁹⁰, O.S. AbouZeid¹⁵⁸, H. Abramowicz¹⁵³, H. Abreu¹⁵², R. Abreu³⁰, Y. Abulaiti^{146a,146b}, B.S. Acharya^{164a,164b,a}, L. Adamczyk^{38a}, D.L. Adams²⁵, J. Adelman¹⁰⁸, S. Adomeit¹⁰⁰, T. Adye¹³¹, A.A. Affolder⁷⁴, T. Agatonovic-Jovin¹³, J.A. Aguilar-Saavedra^{126a,126f}, S.P. Ahlen²², F. Ahmadov^{65,b}, G. Aielli^{133a,133b}, H. Akerstedt^{146a,146b}, T.P.A. Åkesson⁸¹, G. Akimoto¹⁵⁵, A.V. Akimov⁹⁶, G.L. Alberghi^{20a,20b}, J. Albert¹⁶⁹, S. Albrand⁵⁵, M.J. Alconada Verzini⁷¹, M. Aleksa³⁰, I.N. Aleksandrov⁶⁵, C. Alexa^{26a}, G. Alexander¹⁵³, T. Alexopoulos¹⁰, M. Alhroob¹¹³, G. Alimonti^{91a}, L. Alio⁸⁵, J. Alison³¹, S.P. Alkire³⁵, B.M.M. Allbrooke¹⁸, P.P. Allport⁷⁴, A. Aloisio^{104a,104b}, A. Alonso³⁶, F. Alonso⁷¹, C. Alpigiani⁷⁶, A. Altheimer³⁵, B. Alvarez Gonzalez³⁰, D. Álvarez Piqueras¹⁶⁷, M.G. Alvigi^{104a,104b}, B.T. Amadio¹⁵, K. Amako⁶⁶, Y. Amaral Coutinho^{24a}, C. Amelung²³, D. Amidei⁸⁹, S.P. Amor Dos Santos^{126a,126c}, A. Amorim^{126a,126b}, S. Amoroso⁴⁸, N. Amram¹⁵³, G. Amundsen²³, C. Anastopoulos¹³⁹, L.S. Ancu⁴⁹, N. Andari³⁰, T. Andeen³⁵, C.F. Anders^{58b}, G. Anders³⁰, J.K. Anders⁷⁴, K.J. Anderson³¹, A. Andreazza^{91a,91b}, V. Andrei^{58a}, S. Angelidakis⁹, I. Angelozzi¹⁰⁷, P. Anger⁴⁴, A. Angerami³⁵, F. Anghinolfi³⁰, A.V. Anisenkov^{109,c}, N. Anjos¹², A. Annovi^{124a,124b}, M. Antonelli⁴⁷, A. Antonov⁹⁸, J. Antos^{144b}, F. Anulli^{132a}, M. Aoki⁶⁶, L. Aperio Bella¹⁸, G. Arabidze⁹⁰, Y. Arai⁶⁶, J.P. Araque^{126a}, A.T.H. Arce⁴⁵, F.A. Arduh⁷¹, J-F. Arguin⁹⁵, S. Argyropoulos⁴², M. Arik^{19a}, A.J. Armbruster³⁰, O. Arnaez³⁰, V. Arnal⁸², H. Arnold⁴⁸, M. Arratia²⁸, O. Arslan²¹, A. Artamonov⁹⁷, G. Artoni²³, S. Asai¹⁵⁵, N. Asbah⁴², A. Ashkenazi¹⁵³, B. Åsman^{146a,146b}, L. Asquith¹⁴⁹, K. Assamagan²⁵, R. Astalos^{144a}, M. Atkinson¹⁶⁵, N.B. Atlay¹⁴¹, B. Auerbach⁶, K. Augsten¹²⁸, M. Auresseau^{145b}, G. Avolio³⁰, B. Axen¹⁵, M.K. Ayoub¹¹⁷, G. Azuelos^{95,d}, M.A. Baak³⁰, A.E. Baas^{58a}, C. Bacci^{134a,134b}, H. Bachacou¹³⁶, K. Bachas¹⁵⁴, M. Backes³⁰, M. Backhaus³⁰, P. Bagiacchi^{132a,132b}, P. Bagnaia^{132a,132b}, Y. Bai^{33a}, T. Bain³⁵, J.T. Baines¹³¹, O.K. Baker¹⁷⁶, P. Balek¹²⁹, T. Balestri¹⁴⁸, F. Balli⁸⁴, E. Banas³⁹, Sw. Banerjee¹⁷³, A.A.E. Bannoura¹⁷⁵, H.S. Bansil¹⁸, L. Barak³⁰, E.L. Barberio⁸⁸, D. Barberis^{50a,50b}, M. Barbero⁸⁵, T. Barillari¹⁰¹, M. Barisonzi^{164a,164b}, T. Barklow¹⁴³, N. Barlow²⁸, S.L. Barnes⁸⁴, B.M. Barnett¹³¹, R.M. Barnett¹⁵, Z. Barnovska⁵, A. Baroncelli^{134a}, G. Barone⁴⁹, A.J. Barr¹²⁰, F. Barreiro⁸², J. Barreiro Guimarães da Costa⁵⁷, R. Bartoldus¹⁴³, A.E. Barton⁷², P. Bartos^{144a}, A. Basalae¹²³, A. Bassalat¹¹⁷, A. Basye¹⁶⁵, R.L. Bates⁵³, S.J. Batista¹⁵⁸, J.R. Batley²⁸, M. Battaglia¹³⁷, M. Bauge^{132a,132b}, F. Bauer¹³⁶, H.S. Bawa^{143,e}, J.B. Beacham¹¹¹, M.D. Beattie⁷², T. Beau⁸⁰, P.H. Beauchemin¹⁶¹, R. Beccherle^{124a,124b}, P. Bechtel²¹, H.P. Beck^{17,f}, K. Becker¹²⁰, M. Becker⁸³, S. Becker¹⁰⁰, M. Beckingham¹⁷⁰, C. Becot¹¹⁷, A.J. Beddall^{19b}, A. Beddall^{19b}, V.A. Bednyakov⁶⁵, C.P. Bee¹⁴⁸, L.J. Beemster¹⁰⁷, T.A. Beermann¹⁷⁵, M. Begel²⁵, J.K. Behr¹²⁰, C. Belanger-Champagne⁸⁷, W.H. Bell⁴⁹, G. Bella¹⁵³, L. Bellagamba^{20a}, A. Bellerive²⁹, M. Bellomo⁸⁶, K. Belotskiy⁹⁸, O. Beltramello³⁰, O. Benary¹⁵³, D. Benchekroun^{135a}, M. Bender¹⁰⁰, K. Bendtz^{146a,146b}, N. Benekos¹⁰, Y. Benhammou¹⁵³, E. Benhar Noccioli⁴⁹, J.A. Benitez Garcia^{159b}, D.P. Benjamin⁴⁵, J.R. Bensinger²³, S. Bentvelsen¹⁰⁷, L. Beresford¹²⁰, M. Beretta⁴⁷, D. Berge¹⁰⁷, E. Bergeas Kuutmann¹⁶⁶, N. Berger⁵, F. Berghaus¹⁶⁹, J. Beringer¹⁵, C. Bernard²², N.R. Bernard⁸⁶, C. Bernius¹¹⁰, F.U. Bernlochner²¹, T. Berry⁷⁷, P. Berta¹²⁹, C. Bertella⁸³, G. Bertoli^{146a,146b}, F. Bertolucci^{124a,124b}, C. Bertsche¹¹³, D. Bertsche¹¹³, M.I. Besana^{91a}, G.J. Besjes¹⁰⁶, O. Bessidskaia Bylund^{146a,146b}, M. Bessner⁴², N. Besson¹³⁶, C. Betancourt⁴⁸, S. Bethke¹⁰¹, A.J. Bevan⁷⁶, W. Bhimji⁴⁶, R.M. Bianchi¹²⁵, L. Bianchini²³, M. Bianco³⁰, O. Biebel¹⁰⁰, S.P. Bieniek⁷⁸, M. Biglietti^{134a}, J. Bilbao De Mendizabal⁴⁹, H. Bilokon⁴⁷, M. Bindi⁵⁴, S. Binet¹¹⁷, A. Bingul^{19b}, C. Bini^{132a,132b}, C.W. Black¹⁵⁰, J.E. Black¹⁴³, K.M. Black²², D. Blackburn¹³⁸, R.E. Blair⁶, J.-B. Blanchard¹³⁶, J.E. Blanco⁷⁷, T. Blazek^{144a}, I. Bloch⁴², C. Blocker²³, W. Blum^{83,*}, U. Blumenschein⁵⁴, G.J. Bobbink¹⁰⁷,

V.S. Bobrovnikov^{109,c}, S.S. Bocchetta⁸¹, A. Bocci⁴⁵, C. Bock¹⁰⁰, M. Boehler⁴⁸, J.A. Bogaerts³⁰, A.G. Bogdanchikov¹⁰⁹, C. Bohm^{146a}, V. Boisvert⁷⁷, T. Bold^{38a}, V. Boldea^{26a}, A.S. Boldyrev⁹⁹, M. Bomben⁸⁰, M. Bona⁷⁶, M. Boonekamp¹³⁶, A. Borisov¹³⁰, G. Borisso⁷², S. Borroni⁴², J. Bortfeldt¹⁰⁰, V. Bortolotto^{60a,60b,60c}, K. Bos¹⁰⁷, D. Boscherini^{20a}, M. Bosman¹², J. Boudreau¹²⁵, J. Bouffard², E.V. Bouhova-Thacker⁷², D. Boumediene³⁴, C. Bourdarios¹¹⁷, N. Bousson¹¹⁴, A. Boveia³⁰, J. Boyd³⁰, I.R. Boyko⁶⁵, I. Bozic¹³, J. Bracinik¹⁸, A. Brandt⁸, G. Brandt⁵⁴, O. Brandt^{58a}, U. Bratzler¹⁵⁶, B. Brau⁸⁶, J.E. Brau¹¹⁶, H.M. Braun^{175,*}, S.F. Brazzale^{164a,164c}, K. Brendlinger¹²², A.J. Brennan⁸⁸, L. Brenner¹⁰⁷, R. Brenner¹⁶⁶, S. Bressler¹⁷², K. Bristow^{145c}, T.M. Bristow⁴⁶, D. Britton⁵³, D. Britzger⁴², F.M. Brochu²⁸, I. Brock²¹, R. Brock⁹⁰, J. Bronner¹⁰¹, G. Brooijmans³⁵, T. Brooks⁷⁷, W.K. Brooks^{32b}, J. Brosamer¹⁵, E. Brost¹¹⁶, J. Brown⁵⁵, P.A. Bruckman de Renstrom³⁹, D. Bruncko^{144b}, R. Bruneliere⁴⁸, A. Bruni^{20a}, G. Bruni^{20a}, M. Bruschi^{20a}, L. Bryngemark⁸¹, T. Buanes¹⁴, Q. Buat¹⁴², P. Buchholz¹⁴¹, A.G. Buckley⁵³, S.I. Buda^{26a}, I.A. Budagov⁶⁵, F. Buehrer⁴⁸, L. Bugge¹¹⁹, M.K. Bugge¹¹⁹, O. Bulekov⁹⁸, D. Bullock⁸, H. Burckhart³⁰, S. Burdin⁷⁴, B. Burghgrave¹⁰⁸, S. Burke¹³¹, I. Burmeister⁴³, E. Busato³⁴, D. Büscher⁴⁸, V. Büscher⁸³, P. Bussey⁵³, J.M. Butler²², A.I. Butt³, C.M. Buttar⁵³, J.M. Butterworth⁷⁸, P. Butti¹⁰⁷, W. Buttinger²⁵, A. Buzatu⁵³, A.R. Buzykaev^{109,c}, S. Cabrera Urbán¹⁶⁷, D. Caforio¹²⁸, V.M. Cairo^{37a,37b}, O. Cakir^{4a}, P. Calafiura¹⁵, A. Calandri¹³⁶, G. Calderini⁸⁰, P. Calfayan¹⁰⁰, L.P. Caloba^{24a}, D. Calvet³⁴, S. Calvet³⁴, R. Camacho Toro³¹, S. Camarda⁴², P. Camarri^{133a,133b}, D. Cameron¹¹⁹, L.M. Caminada¹⁵, R. Caminal Armadans¹², S. Campana³⁰, M. Campanelli⁷⁸, A. Campoverde¹⁴⁸, V. Canale^{104a,104b}, A. Canepa^{159a}, M. Cano Bret⁷⁶, J. Cantero⁸², R. Cantrill^{126a}, T. Cao⁴⁰, M.D.M. Capeans Garrido³⁰, I. Caprini^{26a}, M. Caprini^{26a}, M. Capua^{37a,37b}, R. Caputo⁸³, R. Cardarelli^{133a}, T. Carli³⁰, G. Carlino^{104a}, L. Carminati^{91a,91b}, S. Caron¹⁰⁶, E. Carquin^{32a}, G.D. Carrillo-Montoya⁸, J.R. Carter²⁸, J. Carvalho^{126a,126c}, D. Casadei⁷⁸, M.P. Casado¹², M. Casolino¹², E. Castaneda-Miranda^{145b}, A. Castelli¹⁰⁷, V. Castillo Gimenez¹⁶⁷, N.F. Castro^{126a,g}, P. Catastini⁵⁷, A. Catinaccio³⁰, J.R. Catmore¹¹⁹, A. Cattai³⁰, J. Caudron⁸³, V. Cavaliere¹⁶⁵, D. Cavalli^{91a}, M. Cavalli-Sforza¹², V. Cavasinni^{124a,124b}, F. Ceradini^{134a,134b}, B.C. Cerio⁴⁵, K. Cerny¹²⁹, A.S. Cerqueira^{24b}, A. Cerri¹⁴⁹, L. Cerrito⁷⁶, F. Cerutti¹⁵, M. Cerv³⁰, A. Cervelli¹⁷, S.A. Cetin^{19c}, A. Chafaq^{135a}, D. Chakraborty¹⁰⁸, I. Chalupkova¹²⁹, P. Chang¹⁶⁵, B. Chapleau⁸⁷, J.D. Chapman²⁸, D.G. Charlton¹⁸, C.C. Chau¹⁵⁸, C.A. Chavez Barajas¹⁴⁹, S. Cheatham¹⁵², A. Chegwidan⁹⁰, S. Chekanov⁶, S.V. Chekulaev^{159a}, G.A. Chelkov^{65,h}, M.A. Chelstowska⁸⁹, C. Chen⁶⁴, H. Chen²⁵, K. Chen¹⁴⁸, L. Chen^{33d,i}, S. Chen^{33c}, X. Chen^{33f}, Y. Chen⁶⁷, H.C. Cheng⁸⁹, Y. Cheng³¹, A. Cheplakov⁶⁵, E. Cheremushkina¹³⁰, R. Cherkaoui El Moursli^{135e}, V. Chernyatin^{25,*}, E. Cheu⁷, L. Chevalier¹³⁶, V. Chiarella⁴⁷, J.T. Childers⁶, G. Chiodini^{73a}, A.S. Chisholm¹⁸, R.T. Chislett⁷⁸, A. Chitan^{26a}, M.V. Chizhov⁶⁵, K. Choi⁶¹, S. Chouridou⁹, B.K.B. Chow¹⁰⁰, V. Christodoulou⁷⁸, D. Chromek-Burckhart³⁰, M.L. Chu¹⁵¹, J. Chudoba¹²⁷, A.J. Chuinard⁸⁷, J.J. Chwastowski³⁹, L. Chytka¹¹⁵, G. Ciapetti^{132a,132b}, A.K. Ciftci^{4a}, D. Cinca⁵³, V. Cindro⁷⁵, I.A. Cioara²¹, A. Ciochio¹⁵, Z.H. Citron¹⁷², M. Ciubancan^{26a}, A. Clark⁴⁹, B.L. Clark⁵⁷, P.J. Clark⁴⁶, R.N. Clarke¹⁵, W. Cleland¹²⁵, C. Clement^{146a,146b}, Y. Coadou⁸⁵, M. Cokal^{164a,164c}, A. Coccaro¹³⁸, J. Cochran⁶⁴, L. Coffey²³, J.G. Cogan¹⁴³, B. Cole³⁵, S. Cole¹⁰⁸, A.P. Colijn¹⁰⁷, J. Collot⁵⁵, T. Colombo^{58c}, G. Compostella¹⁰¹, P. Conde Muiño^{126a,126b}, E. Coniavitis⁴⁸, S.H. Connell^{145b}, I.A. Connelly⁷⁷, S.M. Consonni^{91a,91b}, V. Consorti⁴⁸, S. Constantinescu^{26a}, C. Conta^{121a,121b}, G. Conti³⁰, F. Conventi^{104a,j}, M. Cooke¹⁵, B.D. Cooper⁷⁸, A.M. Cooper-Sarkar¹²⁰, T. Cornelissen¹⁷⁵, M. Corradi^{20a}, F. Corriveau^{87,k}, A. Corso-Radu¹⁶³, A. Cortes-Gonzalez¹², G. Cortiana¹⁰¹, G. Costa^{91a}, M.J. Costa¹⁶⁷, D. Costanzo¹³⁹, D. Côté⁸, G. Cottin²⁸, G. Cowan⁷⁷, B.E. Cox⁸⁴, K. Cranmer¹¹⁰, G. Cree²⁹, S. Crépe-Renaudin⁵⁵, F. Crescioli⁸⁰, W.A. Cribbs^{146a,146b}, M. Crispin Ortuzar¹²⁰, M. Cristinziani²¹, V. Croft¹⁰⁶,

G. Crosetti^{37a,37b}, T. Cuhadar Donszelmann¹³⁹, J. Cummings¹⁷⁶, M. Curatolo⁴⁷, C. Cuthbert¹⁵⁰, H. Czirr¹⁴¹, P. Czodrowski³, S. D'Auria⁵³, M. D'Onofrio⁷⁴, M.J. Da Cunha Sargedas De Sousa^{126a,126b}, C. Da Via⁸⁴, W. Dabrowski^{38a}, A. Dafinca¹²⁰, T. Dai⁸⁹, O. Dale¹⁴, F. Dallaire⁹⁵, C. Dallapiccola⁸⁶, M. Dam³⁶, J.R. Dandoy³¹, N.P. Dang⁴⁸, A.C. Daniells¹⁸, M. Danninger¹⁶⁸, M. Dano Hoffmann¹³⁶, V. Dao⁴⁸, G. Darbo^{50a}, S. Darmora⁸, J. Dassoulas³, A. Dattagupta⁶¹, W. Davey²¹, C. David¹⁶⁹, T. Davidek¹²⁹, E. Davies^{120,l}, M. Davies¹⁵³, P. Davison⁷⁸, Y. Davygora^{58a}, E. Dawe⁸⁸, I. Dawson¹³⁹, R.K. Daya-Ishmukhametova⁸⁶, K. De⁸, R. de Asmundis^{104a}, S. De Castro^{20a,20b}, S. De Cecco⁸⁰, N. De Groot¹⁰⁶, P. de Jong¹⁰⁷, H. De la Torre⁸², F. De Lorenzi⁶⁴, L. De Nooij¹⁰⁷, D. De Pedis^{132a}, A. De Salvo^{132a}, U. De Sanctis¹⁴⁹, A. De Santo¹⁴⁹, J.B. De Vivie De Regie¹¹⁷, W.J. Dearnaley⁷², R. Debbe²⁵, C. Debenedetti¹³⁷, D.V. Dedovich⁶⁵, I. Deigaard¹⁰⁷, J. Del Peso⁸², T. Del Prete^{124a,124b}, D. Delgove¹¹⁷, F. Deliot¹³⁶, C.M. Delitzsch⁴⁹, M. Deliyergiyev⁷⁵, A. Dell'Acqua³⁰, L. Dell'Asta²², M. Dell'Orso^{124a,124b}, M. Della Pietra^{104a,j}, D. della Volpe⁴⁹, M. Delmastro⁵, P.A. Delsart⁵⁵, C. Deluca¹⁰⁷, D.A. DeMarco¹⁵⁸, S. Demers¹⁷⁶, M. Demichev⁶⁵, A. Demilly⁸⁰, S.P. Denisov¹³⁰, D. Derendarz³⁹, J.E. Derkaoui^{135d}, F. Derue⁸⁰, P. Dervan⁷⁴, K. Desch²¹, C. Deterre⁴², P.O. Deviveiros³⁰, A. Dewhurst¹³¹, S. Dhaliwal²³, A. Di Ciaccio^{133a,133b}, L. Di Ciaccio⁵, A. Di Domenico^{132a,132b}, C. Di Donato^{104a,104b}, A. Di Girolamo³⁰, B. Di Girolamo³⁰, A. Di Mattia¹⁵², B. Di Micco^{134a,134b}, R. Di Nardo⁴⁷, A. Di Simone⁴⁸, R. Di Sipio¹⁵⁸, D. Di Valentino²⁹, C. Diaconu⁸⁵, M. Diamond¹⁵⁸, F.A. Dias⁴⁶, M.A. Diaz^{32a}, E.B. Diehl⁸⁹, J. Dietrich¹⁶, S. Diglio⁸⁵, A. Dimitrievska¹³, J. Dingfelder²¹, P. Dita^{26a}, S. Dita^{26a}, F. Dittus³⁰, F. Djama⁸⁵, T. Djobava^{51b}, J.I. Djuvsland^{58a}, M.A.B. do Vale^{24c}, D. Dobos³⁰, M. Dobre^{26a}, C. Doglioni⁴⁹, T. Dohmae¹⁵⁵, J. Dolejsi¹²⁹, Z. Dolezal¹²⁹, B.A. Dolgoshein^{98,*}, M. Donadelli^{24d}, S. Donati^{124a,124b}, P. Dondero^{121a,121b}, J. Donini³⁴, J. Dopke¹³¹, A. Doria^{104a}, M.T. Dova⁷¹, A.T. Doyle⁵³, E. Drechsler⁵⁴, M. Dris¹⁰, E. Dubreuil³⁴, E. Duchovni¹⁷², G. Duckeck¹⁰⁰, O.A. Ducu^{26a,85}, D. Duda¹⁷⁵, A. Dudarev³⁰, L. Dufлот¹¹⁷, L. Duguid⁷⁷, M. Dührssen³⁰, M. Dunford^{58a}, H. Duran Yildiz^{4a}, M. Düren⁵², A. Durglishvili^{51b}, D. Duschinger⁴⁴, M. Dyndal^{38a}, C. Eckardt⁴², K.M. Ecker¹⁰¹, R.C. Edgar⁸⁹, W. Edson², N.C. Edwards⁴⁶, W. Ehrenfeld²¹, T. Eifert³⁰, G. Eigen¹⁴, K. Einsweiler¹⁵, T. Ekelof¹⁶⁶, M. El Kacimi^{135c}, M. Ellert¹⁶⁶, S. Elles⁵, F. Ellinghaus⁸³, A.A. Elliot¹⁶⁹, N. Ellis³⁰, J. Elmsheuser¹⁰⁰, M. Elsing³⁰, D. Emelianov¹³¹, Y. Enari¹⁵⁵, O.C. Endner⁸³, M. Endo¹¹⁸, J. Erdmann⁴³, A. Ereditato¹⁷, G. Ernis¹⁷⁵, J. Ernst², M. Ernst²⁵, S. Errede¹⁶⁵, E. Ertel⁸³, M. Escalier¹¹⁷, H. Esch⁴³, C. Escobar¹²⁵, B. Esposito⁴⁷, A.I. Etienvre¹³⁶, E. Etzion¹⁵³, H. Evans⁶¹, A. Ezhilov¹²³, L. Fabbri^{20a,20b}, G. Facini³¹, R.M. Fakhruddinov¹³⁰, S. Falciano^{132a}, R.J. Falla⁷⁸, J. Faltova¹²⁹, Y. Fang^{33a}, M. Fanti^{91a,91b}, A. Farbin⁸, A. Farilla^{134a}, T. Farooque¹², S. Farrell¹⁵, S.M. Farrington¹⁷⁰, P. Farthouat³⁰, F. Fassi^{135e}, P. Fassnacht³⁰, D. Fassouliotis⁹, M. Fauci Giannelli⁷⁷, A. Favareto^{50a,50b}, L. Fayard¹¹⁷, P. Federic^{144a}, O.L. Fedin^{123,m}, W. Fedorko¹⁶⁸, S. Feigl³⁰, L. Feligioni⁸⁵, C. Feng^{33d}, E.J. Feng⁶, H. Feng⁸⁹, A.B. Fenyuk¹³⁰, P. Fernandez Martinez¹⁶⁷, S. Fernandez Perez³⁰, J. Ferrando⁵³, A. Ferrari¹⁶⁶, P. Ferrari¹⁰⁷, R. Ferrari^{121a}, D.E. Ferreira de Lima⁵³, A. Ferrer¹⁶⁷, D. Ferrere⁴⁹, C. Ferretti⁸⁹, A. Ferretto Parodi^{50a,50b}, M. Fiascaris³¹, F. Fiedler⁸³, A. Filipčić⁷⁵, M. Filipuzzi⁴², F. Filthaut¹⁰⁶, M. Fincke-Keeler¹⁶⁹, K.D. Finelli¹⁵⁰, M.C.N. Fiolhais^{126a,126c}, L. Fiorini¹⁶⁷, A. Firan⁴⁰, A. Fischer², C. Fischer¹², J. Fischer¹⁷⁵, W.C. Fisher⁹⁰, E.A. Fitzgerald²³, M. Flechl⁴⁸, I. Fleck¹⁴¹, P. Fleischmann⁸⁹, S. Fleischmann¹⁷⁵, G.T. Fletcher¹³⁹, G. Fletcher⁷⁶, T. Flick¹⁷⁵, A. Floderus⁸¹, L.R. Flores Castillo^{60a}, M.J. Flowerdew¹⁰¹, A. Formica¹³⁶, A. Forti⁸⁴, D. Fournier¹¹⁷, H. Fox⁷², S. Fracchia¹², P. Francavilla⁸⁰, M. Franchini^{20a,20b}, D. Francis³⁰, L. Franconi¹¹⁹, M. Franklin⁵⁷, M. Fraternali^{121a,121b}, D. Freeborn⁷⁸, S.T. French²⁸, F. Friedrich⁴⁴, D. Froidevaux³⁰, J.A. Frost¹²⁰, C. Fukunaga¹⁵⁶, E. Fullana Torregrosa⁸³, B.G. Fulsom¹⁴³, J. Fuster¹⁶⁷, C. Gabaldon⁵⁵, O. Gabizon¹⁷⁵, A. Gabrielli^{20a,20b}, A. Gabrielli^{132a,132b}, S. Gadatsch¹⁰⁷,

S. Gadomski⁴⁹, G. Gagliardi^{50a,50b}, P. Gagnon⁶¹, C. Galea¹⁰⁶, B. Galhardo^{126a,126c}, E.J. Gallas¹²⁰, B.J. Gallop¹³¹, P. Gallus¹²⁸, G. Galster³⁶, K.K. Gan¹¹¹, J. Gao^{33b,85}, Y. Gao⁴⁶, Y.S. Gao^{143,e}, F.M. Garay Walls⁴⁶, F. Garberson¹⁷⁶, C. García¹⁶⁷, J.E. García Navarro¹⁶⁷, M. Garcia-Sciveres¹⁵, R.W. Gardner³¹, N. Garelli¹⁴³, V. Garonne¹¹⁹, C. Gatti⁴⁷, A. Gaudiello^{50a,50b}, G. Gaudio^{121a}, B. Gaur¹⁴¹, L. Gauthier⁹⁵, P. Gauzzi^{132a,132b}, I.L. Gavrilenko⁹⁶, C. Gay¹⁶⁸, G. Gaycken²¹, E.N. Gazis¹⁰, P. Ge^{33d}, Z. Gecse¹⁶⁸, C.N.P. Gee¹³¹, D.A.A. Geerts¹⁰⁷, Ch. Geich-Gimbel²¹, M.P. Geisler^{58a}, C. Gemme^{50a}, M.H. Genest⁵⁵, S. Gentile^{132a,132b}, M. George⁵⁴, S. George⁷⁷, D. Gerbaudo¹⁶³, A. Gershon¹⁵³, H. Ghazlane^{135b}, B. Giacobbe^{20a}, S. Giagu^{132a,132b}, V. Giangiobbe¹², P. Giannetti^{124a,124b}, B. Gibbard²⁵, S.M. Gibson⁷⁷, M. Gilchriese¹⁵, T.P.S. Gillam²⁸, D. Gillberg³⁰, G. Gilles³⁴, D.M. Gingrich^{3,d}, N. Giokaris⁹, M.P. Giordani^{164a,164c}, F.M. Giorgi^{20a}, F.M. Giorgi¹⁶, P.F. Giraud¹³⁶, P. Giromini⁴⁷, D. Giugni^{91a}, C. Giuliani⁴⁸, M. Giuliani^{58b}, B.K. Gjølsten¹¹⁹, S. Gkaitatzis¹⁵⁴, I. Gkialas¹⁵⁴, E.L. Gkougkousis¹¹⁷, L.K. Gladilin⁹⁹, C. Glasman⁸², J. Glatzer³⁰, P.C.F. Glaysheer⁴⁶, A. Glazov⁴², M. Goblirsch-Kolb¹⁰¹, J.R. Goddard⁷⁶, J. Godlewski³⁹, S. Goldfarb⁸⁹, T. Golling⁴⁹, D. Golubkov¹³⁰, A. Gomes^{126a,126b,126d}, R. Gonçalo^{126a}, J. Goncalves Pinto Firmino Da Costa¹³⁶, L. Gonella²¹, S. González de la Hoz¹⁶⁷, G. Gonzalez Parra¹², S. Gonzalez-Sevilla⁴⁹, L. Goossens³⁰, P.A. Gorbounov⁹⁷, H.A. Gordon²⁵, I. Gorelov¹⁰⁵, B. Gorini³⁰, E. Gorini^{73a,73b}, A. Gorišek⁷⁵, E. Gornicki³⁹, A.T. Goshaw⁴⁵, C. Gössling⁴³, M.I. Gostkin⁶⁵, D. Goudami^{135c}, A.G. Goussiou¹³⁸, N. Govender^{145b}, H.M.X. Grabas¹³⁷, L. Graber⁵⁴, I. Grabowska-Bold^{38a}, P. Grafström^{20a,20b}, K.-J. Grahm⁴², J. Gramling⁴⁹, E. Gramstad¹¹⁹, S. Grancagnolo¹⁶, V. Grassi¹⁴⁸, V. Gratchev¹²³, H.M. Gray³⁰, E. Graziani^{134a}, Z.D. Greenwood^{79,n}, K. Gregersen⁷⁸, I.M. Gregor⁴², P. Grenier¹⁴³, J. Griffiths⁸, A.A. Grillo¹³⁷, K. Grimm⁷², S. Grinstein^{12,o}, Ph. Gris³⁴, J.-F. Grivaz¹¹⁷, J.P. Grohs⁴⁴, A. Grohsjean⁴², E. Gross¹⁷², J. Grosse-Knetter⁵⁴, G.C. Grossi⁷⁹, Z.J. Grout¹⁴⁹, L. Guan^{33b}, J. Guenther¹²⁸, F. Guescini⁴⁹, D. Guest¹⁷⁶, O. Gueta¹⁵³, E. Guido^{50a,50b}, T. Guillemin¹¹⁷, S. Guindon², U. Gul⁵³, C. Gumpert⁴⁴, J. Guo^{33e}, S. Gupta¹²⁰, P. Gutierrez¹¹³, N.G. Gutierrez Ortiz⁵³, C. Gutsche⁴⁴, C. Guyot¹³⁶, C. Gwenlan¹²⁰, C.B. Gwilliam⁷⁴, A. Haas¹¹⁰, C. Haber¹⁵, H.K. Hadavand⁸, N. Haddad^{135e}, P. Haefner²¹, S. Hageböck²¹, Z. Hajduk³⁹, H. Hakobyan¹⁷⁷, M. Haleem⁴², J. Haley¹¹⁴, D. Hall¹²⁰, G. Halladjian⁹⁰, G.D. Hallewell⁸⁵, K. Hamacher¹⁷⁵, P. Hamal¹¹⁵, K. Hamano¹⁶⁹, M. Hamer⁵⁴, A. Hamilton^{145a}, G.N. Hamity^{145c}, P.G. Hamnett⁴², L. Han^{33b}, K. Hanagaki¹¹⁸, K. Hanawa¹⁵⁵, M. Hance¹⁵, P. Hanke^{58a}, R. Hanna¹³⁶, J.B. Hansen³⁶, J.D. Hansen³⁶, M.C. Hansen²¹, P.H. Hansen³⁶, K. Hara¹⁶⁰, A.S. Hard¹⁷³, T. Harenberg¹⁷⁵, F. Hariri¹¹⁷, S. Harkusha⁹², R.D. Harrington⁴⁶, P.F. Harrison¹⁷⁰, F. Hartjes¹⁰⁷, M. Hasegawa⁶⁷, S. Hasegawa¹⁰³, Y. Hasegawa¹⁴⁰, A. Hasib¹¹³, S. Hassani¹³⁶, S. Haug¹⁷, R. Hauser⁹⁰, L. Hauswald⁴⁴, M. Havranek¹²⁷, C.M. Hawkes¹⁸, R.J. Hawkings³⁰, A.D. Hawkins⁸¹, T. Hayashi¹⁶⁰, D. Hayden⁹⁰, C.P. Hays¹²⁰, J.M. Hays⁷⁶, H.S. Hayward⁷⁴, S.J. Haywood¹³¹, S.J. Head¹⁸, T. Heck⁸³, V. Hedberg⁸¹, L. Heelan⁸, S. Heim¹²², T. Heim¹⁷⁵, B. Heinemann¹⁵, L. Heinrich¹¹⁰, J. Hejbal¹²⁷, L. Helary²², S. Hellman^{146a,146b}, D. Hellmich²¹, C. Helsens³⁰, J. Henderson¹²⁰, R.C.W. Henderson⁷², Y. Heng¹⁷³, C. Hengler⁴², A. Henrichs¹⁷⁶, A.M. Henriques Correia³⁰, S. Henrot-Versille¹¹⁷, G.H. Herbert¹⁶, Y. Hernández Jiménez¹⁶⁷, R. Herrberg-Schubert¹⁶, G. Herten⁴⁸, R. Hertenberger¹⁰⁰, L. Hervas³⁰, G.G. Hesketh⁷⁸, N.P. Hessey¹⁰⁷, J.W. Hetherly⁴⁰, R. Hickling⁷⁶, E. Higón-Rodríguez¹⁶⁷, E. Hill¹⁶⁹, J.C. Hill²⁸, K.H. Hiller⁴², S.J. Hillier¹⁸, I. Hinchliffe¹⁵, E. Hines¹²², R.R. Hinman¹⁵, M. Hirose¹⁵⁷, D. Hirschbuehl¹⁷⁵, J. Hobbs¹⁴⁸, N. Hod¹⁰⁷, M.C. Hodgkinson¹³⁹, P. Hodgson¹³⁹, A. Hoecker³⁰, M.R. Hoefkamp¹⁰⁵, F. Hoenig¹⁰⁰, M. Hohlfield⁸³, D. Hohn²¹, T.R. Holmes¹⁵, M. Homann⁴³, T.M. Hong¹²⁵, L. Hooft van Huysduynen¹¹⁰, W.H. Hopkins¹¹⁶, Y. Horii¹⁰³, A.J. Horton¹⁴², J.-Y. Hostachy⁵⁵, S. Hou¹⁵¹, A. Hoummada^{135a}, J. Howard¹²⁰, J. Howarth⁴², M. Hrabovsky¹¹⁵, I. Hristova¹⁶, J. Hrivnac¹¹⁷, T. Hryn'ova⁵, A. Hrynevich⁹³, C. Hsu^{145c}, P.J. Hsu^{151,p}, S.-C. Hsu¹³⁸, D. Hu³⁵, Q. Hu^{33b}, X. Hu⁸⁹, Y. Huang⁴², Z. Hubacek³⁰, F. Hubaut⁸⁵,

F. Huegging²¹, T.B. Huffman¹²⁰, E.W. Hughes³⁵, G. Hughes⁷², M. Huhtinen³⁰, T.A. Hülsing⁸³, N. Huseynov^{65,b}, J. Huston⁹⁰, J. Huth⁵⁷, G. Iacobucci⁴⁹, G. Iakovidis²⁵, I. Ibragimov¹⁴¹, L. Iconomidou-Fayard¹¹⁷, E. Ideal¹⁷⁶, Z. Idrissi^{135e}, P. Iengo³⁰, O. Igonkina¹⁰⁷, T. Iizawa¹⁷¹, Y. Ikegami⁶⁶, K. Ikematsu¹⁴¹, M. Ikeno⁶⁶, Y. Ilchenko^{31,q}, D. Iliadis¹⁵⁴, N. Ilic¹⁴³, Y. Inamaru⁶⁷, T. Ince¹⁰¹, P. Ioannou⁹, M. Iodice^{134a}, K. Iordanidou³⁵, V. Ippolito⁵⁷, A. Irles Quiles¹⁶⁷, C. Isaksson¹⁶⁶, M. Ishino⁶⁸, M. Ishitsuka¹⁵⁷, R. Ishmukhametov¹¹¹, C. Issever¹²⁰, S. Istin^{19a}, J.M. Iturbe Ponce⁸⁴, R. Iuppa^{133a,133b}, J. Ivarsson⁸¹, W. Iwanski³⁹, H. Iwasaki⁶⁶, J.M. Izen⁴¹, V. Izzo^{104a}, S. Jabbar³, B. Jackson¹²², M. Jackson⁷⁴, P. Jackson¹, M.R. Jaekel³⁰, V. Jain², K. Jakobs⁴⁸, S. Jakobsen³⁰, T. Jakoubek¹²⁷, J. Jakubek¹²⁸, D.O. Jamin¹⁵¹, D.K. Jana⁷⁹, E. Jansen⁷⁸, R. Jansky⁶², J. Janssen²¹, M. Janus¹⁷⁰, G. Jarlskog⁸¹, N. Javadov^{65,b}, T. Javůrek⁴⁸, L. Jeanty¹⁵, J. Jejelava^{51a,r}, G.-Y. Jeng¹⁵⁰, D. Jennens⁸⁸, P. Jenni^{48,s}, J. Jentzsch⁴³, C. Jeske¹⁷⁰, S. Jézéquel⁵, H. Ji¹⁷³, J. Jia¹⁴⁸, Y. Jiang^{33b}, S. Jiggins⁷⁸, J. Jimenez Pena¹⁶⁷, S. Jin^{33a}, A. Jinaru^{26a}, O. Jinnouchi¹⁵⁷, M.D. Joergensen³⁶, P. Johansson¹³⁹, K.A. Johns⁷, K. Jon-And^{146a,146b}, G. Jones¹⁷⁰, R.W.L. Jones⁷², T.J. Jones⁷⁴, J. Jongmanns^{58a}, P.M. Jorge^{126a,126b}, K.D. Joshi⁸⁴, J. Jovicevic^{159a}, X. Ju¹⁷³, C.A. Jung⁴³, P. Jussel⁶², A. Juste Rozas^{12,o}, M. Kaci¹⁶⁷, A. Kaczmarska³⁹, M. Kado¹¹⁷, H. Kagan¹¹¹, M. Kagan¹⁴³, S.J. Kahn⁸⁵, E. Kajomovitz⁴⁵, C.W. Kalderon¹²⁰, S. Kama⁴⁰, A. Kamenshchikov¹³⁰, N. Kanaya¹⁵⁵, M. Kaneda³⁰, S. Kaneti²⁸, V.A. Kantserov⁹⁸, J. Kanzaki⁶⁶, B. Kaplan¹¹⁰, A. Kapliy³¹, D. Kar⁵³, K. Karakostas¹⁰, A. Karamaoun³, N. Karastathis^{10,107}, M.J. Kareem⁵⁴, M. Karnevskiy⁸³, S.N. Karpov⁶⁵, Z.M. Karpova⁶⁵, K. Karthik¹¹⁰, V. Kartvelishvili⁷², A.N. Karyukhin¹³⁰, L. Kashif¹⁷³, R.D. Kass¹¹¹, A. Kastanas¹⁴, Y. Kataoka¹⁵⁵, A. Katre⁴⁹, J. Katzy⁴², K. Kawagoe⁷⁰, T. Kawamoto¹⁵⁵, G. Kawamura⁵⁴, S. Kazama¹⁵⁵, V.F. Kazanin^{109,c}, M.Y. Kazarinov⁶⁵, R. Keeler¹⁶⁹, R. Kehoe⁴⁰, J.S. Keller⁴², J.J. Kempster⁷⁷, H. Keoshkerian⁸⁴, O. Kepka¹²⁷, B.P. Kerševan⁷⁵, S. Kersten¹⁷⁵, R.A. Keyes⁸⁷, F. Khalil-zada¹¹, H. Khandanyan^{146a,146b}, A. Khanov¹¹⁴, A.G. Kharlamov^{109,c}, T.J. Khoo²⁸, V. Khovanskiy⁹⁷, E. Khramov⁶⁵, J. Khubua^{51b,t}, H.Y. Kim⁸, H. Kim^{146a,146b}, S.H. Kim¹⁶⁰, Y. Kim³¹, N. Kimura¹⁵⁴, O.M. Kind¹⁶, B.T. King⁷⁴, M. King¹⁶⁷, R.S.B. King¹²⁰, S.B. King¹⁶⁸, J. Kirk¹³¹, A.E. Kiryunin¹⁰¹, T. Kishimoto⁶⁷, D. Kisielewska^{38a}, F. Kiss⁴⁸, K. Kiuchi¹⁶⁰, O. Kivernyk¹³⁶, E. Kladiva^{144b}, M.H. Klein³⁵, M. Klein⁷⁴, U. Klein⁷⁴, K. Kleinknecht⁸³, P. Klimek^{146a,146b}, A. Klimentov²⁵, R. Klingenberg⁴³, J.A. Klinger⁸⁴, T. Klioutchnikova³⁰, E.-E. Kluge^{58a}, P. Kluit¹⁰⁷, S. Kluth¹⁰¹, E. Kneringer⁶², E.B.F.G. Knoops⁸⁵, A. Knue⁵³, A. Kobayashi¹⁵⁵, D. Kobayashi¹⁵⁷, T. Kobayashi¹⁵⁵, M. Kobel⁴⁴, M. Kocian¹⁴³, P. Kodys¹²⁹, T. Koffas²⁹, E. Koffeman¹⁰⁷, L.A. Kogan¹²⁰, S. Kohlmann¹⁷⁵, Z. Kohout¹²⁸, T. Kohriki⁶⁶, T. Koi¹⁴³, H. Kolanoski¹⁶, I. Koletsou⁵, A.A. Komar^{96,*}, Y. Komori¹⁵⁵, T. Kondo⁶⁶, N. Kondrashova⁴², K. Köneke⁴⁸, A.C. König¹⁰⁶, S. König⁸³, T. Kono^{66,u}, R. Konoplich^{110,v}, N. Konstantinidis⁷⁸, R. Kopeliansky¹⁵², S. Koperny^{38a}, L. Köpke⁸³, A.K. Kopp⁴⁸, K. Korcyl³⁹, K. Kordas¹⁵⁴, A. Korn⁷⁸, A.A. Korol^{109,c}, I. Korolkov¹², E.V. Korolkova¹³⁹, O. Kortner¹⁰¹, S. Kortner¹⁰¹, T. Kosek¹²⁹, V.V. Kostyukhin²¹, V.M. Kotov⁶⁵, A. Kotwal⁴⁵, A. Kourkoumeli-Charalampidi¹⁵⁴, C. Kourkoumelis⁹, V. Kouskoura²⁵, A. Koutsman^{159a}, R. Kowalewski¹⁶⁹, T.Z. Kowalski^{38a}, W. Kozanecki¹³⁶, A.S. Kozhin¹³⁰, V.A. Kramarenko⁹⁹, G. Kramberger⁷⁵, D. Krasnopevtsev⁹⁸, M.W. Krasny⁸⁰, A. Krasznahorkay³⁰, J.K. Kraus²¹, A. Kravchenko²⁵, S. Kreiss¹¹⁰, M. Kretz^{58c}, J. Kretzschmar⁷⁴, K. Kreutzfeldt⁵², P. Krieger¹⁵⁸, K. Krizka³¹, K. Kroeninger⁴³, H. Kroha¹⁰¹, J. Kroll¹²², J. Kruseberg²¹, J. Krstic¹³, U. Kruchonak⁶⁵, H. Krüger²¹, N. Krumnack⁶⁴, Z.V. Krumshcheyn⁶⁵, A. Kruse¹⁷³, M.C. Kruse⁴⁵, M. Kruskal²², T. Kubota⁸⁸, H. Kucuk⁷⁸, S. Kудay^{4b}, S. Kuehn⁴⁸, A. Kugel^{58c}, F. Kuger¹⁷⁴, A. Kuhl¹³⁷, T. Kuhl⁴², V. Kukhtin⁶⁵, Y. Kulchitsky⁹², S. Kuleshov^{32b}, M. Kuna^{132a,132b}, T. Kunigo⁶⁸, A. Kupco¹²⁷, H. Kurashige⁶⁷, Y.A. Kurochkin⁹², R. Kurumida⁶⁷, V. Kus¹²⁷, E.S. Kuwertz¹⁶⁹, M. Kuze¹⁵⁷, J. Kvita¹¹⁵, T. Kwan¹⁶⁹, D. Kyriazopoulos¹³⁹, A. La Rosa⁴⁹, J.L. La Rosa Navarro^{24d}, L. La Rotonda^{37a,37b},

C. Lacasta¹⁶⁷, F. Lacava^{132a,132b}, J. Lacey²⁹, H. Lacker¹⁶, D. Lacour⁸⁰, V.R. Lacuesta¹⁶⁷, E. Ladygin⁶⁵, R. Lafaye⁵, B. Laforge⁸⁰, T. Lagouri¹⁷⁶, S. Lai⁴⁸, L. Lambourne⁷⁸, S. Lammers⁶¹, C.L. Lampen⁷, W. Lampl⁷, E. Lançon¹³⁶, U. Landgraf⁴⁸, M.P.J. Landon⁷⁶, V.S. Lang^{58a}, J.C. Lange¹², A.J. Lankford¹⁶³, F. Lanni²⁵, K. Lantzsch³⁰, S. Laplace⁸⁰, C. Lapoire³⁰, J.F. Laporte¹³⁶, T. Lari^{91a}, F. Lasagni Manghi^{20a,20b}, M. Lassnig³⁰, P. Laurelli⁴⁷, W. Lavrijsen¹⁵, A.T. Law¹³⁷, P. Laycock⁷⁴, T. Lazovich⁵⁷, O. Le Dortz⁸⁰, E. Le Guirriec⁸⁵, E. Le Menedeu¹², M. LeBlanc¹⁶⁹, T. LeCompte⁶, F. Ledroit-Guillon⁵⁵, C.A. Lee^{145b}, S.C. Lee¹⁵¹, L. Lee¹, G. Lefebvre⁸⁰, M. Lefebvre¹⁶⁹, F. Legger¹⁰⁰, C. Leggett¹⁵, A. Lehan⁷⁴, G. Lehmann Miotto³⁰, X. Lei⁷, W.A. Leight²⁹, A. Leisos^{154,w}, A.G. Leister¹⁷⁶, M.A.L. Leite^{24d}, R. Leitner¹²⁹, D. Lellouch¹⁷², B. Lemmer⁵⁴, K.J.C. Leney⁷⁸, T. Lenz²¹, B. Lenzi³⁰, R. Leone⁷, S. Leone^{124a,124b}, C. Leonidopoulos⁴⁶, S. Leontsinis¹⁰, C. Leroy⁹⁵, C.G. Lester²⁸, M. Levchenko¹²³, J. Levêque⁵, D. Levin⁸⁹, L.J. Levinson¹⁷², M. Levy¹⁸, A. Lewis¹²⁰, A.M. Leyko²¹, M. Leyton⁴¹, B. Li^{33b,x}, H. Li¹⁴⁸, H.L. Li³¹, L. Li⁴⁵, L. Li^{33e}, S. Li⁴⁵, Y. Li^{33c,y}, Z. Liang¹³⁷, H. Liao³⁴, B. Liberti^{133a}, A. Liblong¹⁵⁸, P. Lichard³⁰, K. Lie¹⁶⁵, J. Liebal²¹, W. Liebig¹⁴, C. Limbach²¹, A. Limosani¹⁵⁰, S.C. Lin^{151,z}, T.H. Lin⁸³, F. Linde¹⁰⁷, B.E. Lindquist¹⁴⁸, J.T. Linnemann⁹⁰, E. Lipeles¹²², A. Lipniacka¹⁴, M. Lisovyi^{58b}, T.M. Liss¹⁶⁵, D. Lissauer²⁵, A. Lister¹⁶⁸, A.M. Litke¹³⁷, B. Liu^{151,aa}, D. Liu¹⁵¹, J. Liu⁸⁵, J.B. Liu^{33b}, K. Liu⁸⁵, L. Liu¹⁶⁵, M. Liu⁴⁵, M. Liu^{33b}, Y. Liu^{33b}, M. Livan^{121a,121b}, A. Lleres⁵⁵, J. Llorente Merino⁸², S.L. Lloyd⁷⁶, F. Lo Sterzo¹⁵¹, E. Lobodzinska⁴², P. Loch⁷, W.S. Lockman¹³⁷, F.K. Loebinger⁸⁴, A.E. Loevschall-Jensen³⁶, A. Loginov¹⁷⁶, T. Lohse¹⁶, K. Lohwasser⁴², M. Lokajicek¹²⁷, B.A. Long²², J.D. Long⁸⁹, R.E. Long⁷², K.A. Looper¹¹¹, L. Lopes^{126a}, D. Lopez Mateos⁵⁷, B. Lopez Paredes¹³⁹, I. Lopez Paz¹², J. Lorenz¹⁰⁰, N. Lorenzo Martinez⁶¹, M. Losada¹⁶², P. Loscutoff¹⁵, P.J. Lösel¹⁰⁰, X. Lou^{33a}, A. Lounis¹¹⁷, J. Love⁶, P.A. Love⁷², N. Lu⁸⁹, H.J. Lubatti¹³⁸, C. Luci^{132a,132b}, A. Lucotte⁵⁵, F. Luehring⁶¹, W. Lukas⁶², L. Luminari^{132a}, O. Lundberg^{146a,146b}, B. Lund-Jensen¹⁴⁷, D. Lynn²⁵, R. Lysak¹²⁷, E. Lytken⁸¹, H. Ma²⁵, L.L. Ma^{33d}, G. Maccarrone⁴⁷, A. Macchiolo¹⁰¹, C.M. Macdonald¹³⁹, J. Machado Miguens^{122,126b}, D. Macina³⁰, D. Madaffari⁸⁵, R. Madar³⁴, H.J. Maddocks⁷², W.F. Mader⁴⁴, A. Madsen¹⁶⁶, S. Maeland¹⁴, T. Maeno²⁵, A. Maevskiy⁹⁹, E. Magradze⁵⁴, K. Mahboubi⁴⁸, J. Mahlstedt¹⁰⁷, C. Maiani¹³⁶, C. Maidantchik^{24a}, A.A. Maier¹⁰¹, T. Maier¹⁰⁰, A. Maio^{126a,126b,126d}, S. Majewski¹¹⁶, Y. Makida⁶⁶, N. Makovec¹¹⁷, B. Malaescu⁸⁰, Pa. Malecki³⁹, V.P. Maleev¹²³, F. Malek⁵⁵, U. Mallik⁶³, D. Malon⁶, C. Malone¹⁴³, S. Maltezos¹⁰, V.M. Malyshev¹⁰⁹, S. Malyukov³⁰, J. Mamuzic⁴², G. Mancini⁴⁷, B. Mandelli³⁰, L. Mandelli^{91a}, I. Mandić⁷⁵, R. Mandrysch⁶³, J. Maneira^{126a,126b}, A. Manfredini¹⁰¹, L. Manhaes de Andrade Filho^{24b}, J. Manjarres Ramos^{159b}, A. Mann¹⁰⁰, P.M. Manning¹³⁷, A. Manousakis-Katsikakis⁹, B. Mansoulie¹³⁶, R. Mantifel⁸⁷, M. Mantoani⁵⁴, L. Mapelli³⁰, L. March^{145c}, G. Marchiori⁸⁰, M. Marcisovsky¹²⁷, C.P. Marino¹⁶⁹, M. Marjanovic¹³, F. Marroquim^{24a}, S.P. Marsden⁸⁴, Z. Marshall¹⁵, L.F. Marti¹⁷, S. Marti-Garcia¹⁶⁷, B. Martin⁹⁰, T.A. Martin¹⁷⁰, V.J. Martin⁴⁶, B. Martin dit Latour¹⁴, M. Martinez^{12,o}, S. Martin-Haugh¹³¹, V.S. Martoiu^{26a}, A.C. Martyniuk⁷⁸, M. Marx¹³⁸, F. Marzano^{132a}, A. Marzin³⁰, L. Masetti⁸³, T. Mashimo¹⁵⁵, R. Mashinistov⁹⁶, J. Masik⁸⁴, A.L. Maslennikov^{109,c}, I. Massa^{20a,20b}, L. Massa^{20a,20b}, N. Massol⁵, P. Mastrandrea¹⁴⁸, A. Mastroberardino^{37a,37b}, T. Masubuchi¹⁵⁵, P. Mättig¹⁷⁵, J. Mattmann⁸³, J. Maurer^{26a}, S.J. Maxfield⁷⁴, D.A. Maximov^{109,c}, R. Mazini¹⁵¹, S.M. Mazza^{91a,91b}, L. Mazzaferro^{133a,133b}, G. Mc Goldrick¹⁵⁸, S.P. Mc Kee⁸⁹, A. McCarn⁸⁹, R.L. McCarthy¹⁴⁸, T.G. McCarthy²⁹, N.A. McCubbin¹³¹, K.W. McFarlane^{56,*}, J.A. Mcfayden⁷⁸, G. Mchedlidze⁵⁴, S.J. McMahon¹³¹, R.A. McPherson^{169,k}, M. Medinnis⁴², S. Meehan^{145a}, S. Mehlhase¹⁰⁰, A. Mehta⁷⁴, K. Meier^{58a}, C. Meineck¹⁰⁰, B. Meirose⁴¹, B.R. Mellado Garcia^{145c}, F. Meloni¹⁷, A. Mengarelli^{20a,20b}, S. Menke¹⁰¹, E. Meoni¹⁶¹, K.M. Mercurio⁵⁷, S. Mergelmeyer²¹, P. Mermod⁴⁹, L. Merola^{104a,104b}, C. Meroni^{91a}, F.S. Merritt³¹, A. Messina^{132a,132b}, J. Metcalfe²⁵, A.S. Mete¹⁶³, C. Meyer⁸³, C. Meyer¹²², J-P. Meyer¹³⁶, J. Meyer¹⁰⁷, R.P. Middleton¹³¹,

S. Miglioranzì^{164a,164c}, L. Mijović²¹, G. Mikenberg¹⁷², M. Mikestikova¹²⁷, M. Mikuz⁷⁵, M. Milesi⁸⁸, A. Milic³⁰, D.W. Miller³¹, C. Mills⁴⁶, A. Milov¹⁷², D.A. Milstead^{146a,146b}, A.A. Minaenko¹³⁰, Y. Minami¹⁵⁵, I.A. Minashvili⁶⁵, A.I. Mincer¹¹⁰, B. Mindur^{38a}, M. Mineev⁶⁵, Y. Ming¹⁷³, L.M. Mir¹², T. Mitani¹⁷¹, J. Mitrevski¹⁰⁰, V.A. Mitsou¹⁶⁷, A. Miucci⁴⁹, P.S. Miyagawa¹³⁹, J.U. Mjörnmark⁸¹, T. Moa^{146a,146b}, K. Mochizuki⁸⁵, S. Mohapatra³⁵, W. Mohr⁴⁸, S. Molander^{146a,146b}, R. Moles-Valls¹⁶⁷, K. Mönig⁴², C. Monini⁵⁵, J. Monk³⁶, E. Monnier⁸⁵, J. Montejo Berlingen¹², F. Monticelli⁷¹, S. Monzani^{132a,132b}, R.W. Moore³, N. Morange¹¹⁷, D. Moreno¹⁶², M. Moreno Llácer⁵⁴, P. Morettini^{50a}, M. Morgenstern⁴⁴, M. Morii⁵⁷, M. Morinaga¹⁵⁵, V. Morisbak¹¹⁹, S. Moritz⁸³, A.K. Morley¹⁴⁷, G. Mornacchi³⁰, J.D. Morris⁷⁶, S.S. Mortensen³⁶, A. Morton⁵³, L. Morvaj¹⁰³, M. Mosidze^{51b}, J. Moss¹¹¹, K. Motohashi¹⁵⁷, R. Mount¹⁴³, E. Mountricha²⁵, S.V. Mouraviev^{96,*}, E.J.W. Moyse⁸⁶, S. Muanza⁸⁵, R.D. Mudd¹⁸, F. Mueller¹⁰¹, J. Mueller¹²⁵, K. Mueller²¹, R.S.P. Mueller¹⁰⁰, T. Mueller²⁸, D. Muenstermann⁴⁹, P. Mullen⁵³, Y. Munwes¹⁵³, J.A. Murillo Quijada¹⁸, W.J. Murray^{170,131}, H. Musheghyan⁵⁴, E. Musto¹⁵², A.G. Myagkov^{130,ab}, M. Myska¹²⁸, O. Nackenhorst⁵⁴, J. Nadal⁵⁴, K. Nagai¹²⁰, R. Nagai¹⁵⁷, Y. Nagai⁸⁵, K. Nagano⁶⁶, A. Nagarkar¹¹¹, Y. Nagasaka⁵⁹, K. Nagata¹⁶⁰, M. Nagel¹⁰¹, E. Nagy⁸⁵, A.M. Nairz³⁰, Y. Nakahama³⁰, K. Nakamura⁶⁶, T. Nakamura¹⁵⁵, I. Nakano¹¹², H. Namasivayam⁴¹, R.F. Naranjo Garcia⁴², R. Narayan³¹, T. Naumann⁴², G. Navarro¹⁶², R. Nayyar⁷, H.A. Neal⁸⁹, P.Yu. Nechaeva⁹⁶, T.J. Neep⁸⁴, P.D. Nef¹⁴³, A. Negri^{121a,121b}, M. Negrini^{20a}, S. Nektarijevic¹⁰⁶, C. Nellist¹¹⁷, A. Nelson¹⁶³, S. Nemecek¹²⁷, P. Nemethy¹¹⁰, A.A. Nepomuceno^{24a}, M. Nessi^{30,ac}, M.S. Neubauer¹⁶⁵, M. Neumann¹⁷⁵, R.M. Neves¹¹⁰, P. Nevski²⁵, P.R. Newman¹⁸, D.H. Nguyen⁶, R.B. Nickerson¹²⁰, R. Nicolaidou¹³⁶, B. Nicquevert³⁰, J. Nielsen¹³⁷, N. Nikiforou³⁵, A. Nikiforov¹⁶, V. Nikolaenko^{130,ab}, I. Nikolic-Audit⁸⁰, K. Nikolopoulos¹⁸, J.K. Nilsen¹¹⁹, P. Nilsson²⁵, Y. Ninomiya¹⁵⁵, A. Nisati^{132a}, R. Nisius¹⁰¹, T. Nobe¹⁵⁷, M. Nomachi¹¹⁸, I. Nomidis²⁹, T. Nooney⁷⁶, S. Norberg¹¹³, M. Nordberg³⁰, O. Novgorodova⁴⁴, S. Nowak¹⁰¹, M. Nozaki⁶⁶, L. Nozka¹¹⁵, K. Ntekas¹⁰, G. Nunes Hanninger⁸⁸, T. Nunnemann¹⁰⁰, E. Nurse⁷⁸, F. Nuti⁸⁸, B.J. O'Brien⁴⁶, F. O'grady⁷, D.C. O'Neil¹⁴², V. O'Shea⁵³, F.G. Oakham^{29,d}, H. Oberlack¹⁰¹, T. Obermann²¹, J. Ocariz⁸⁰, A. Ochi⁶⁷, I. Ochoa⁷⁸, J.P. Ochoa-Ricoux^{32a}, S. Oda⁷⁰, S. Odaka⁶⁶, H. Ogren⁶¹, A. Oh⁸⁴, S.H. Oh⁴⁵, C.C. Ohm¹⁵, H. Ohman¹⁶⁶, H. Oide³⁰, W. Okamura¹¹⁸, H. Okawa¹⁶⁰, Y. Okumura³¹, T. Okuyama¹⁵⁵, A. Olariu^{26a}, S.A. Olivares Pino⁴⁶, D. Oliveira Damazio²⁵, E. Oliver Garcia¹⁶⁷, A. Olszewski³⁹, J. Olszowska³⁹, A. Onofre^{126a,126e}, P.U.E. Onyisi^{31,q}, C.J. Oram^{159a}, M.J. Oreglia³¹, Y. Oren¹⁵³, D. Orestano^{134a,134b}, N. Orlando¹⁵⁴, C. Oropeza Barrera⁵³, R.S. Orr¹⁵⁸, B. Osculati^{50a,50b}, R. Ospanov⁸⁴, G. Otero y Garzon²⁷, H. Otono⁷⁰, M. Ouchrif^{135d}, E.A. Ouellette¹⁶⁹, F. Ould-Saada¹¹⁹, A. Ouraou¹³⁶, K.P. Oussoren¹⁰⁷, Q. Ouyang^{33a}, A. Ovcharova¹⁵, M. Owen⁵³, R.E. Owen¹⁸, V.E. Ozcan^{19a}, N. Ozturk⁸, K. Pachal¹⁴², A. Pacheco Pages¹², C. Padilla Aranda¹², M. Pagáčová⁴⁸, S. Pagan Griso¹⁵, E. Paganis¹³⁹, C. Pahl¹⁰¹, F. Paige²⁵, P. Pais⁸⁶, K. Pajchel¹¹⁹, G. Palacino^{159b}, S. Palestini³⁰, M. Palka^{38b}, D. Pallin³⁴, A. Palma^{126a,126b}, Y.B. Pan¹⁷³, E. Panagiotopoulou¹⁰, C.E. Pandini⁸⁰, J.G. Panduro Vazquez⁷⁷, P. Pani^{146a,146b}, S. Panitkin²⁵, D. Pantea^{26a}, L. Paolozzi⁴⁹, Th.D. Papadopoulou¹⁰, K. Papageorgiou¹⁵⁴, A. Paramonov⁶, D. Paredes Hernandez¹⁵⁴, M.A. Parker²⁸, K.A. Parker¹³⁹, F. Parodi^{50a,50b}, J.A. Parsons³⁵, U. Parzefall⁴⁸, E. Pasqualucci^{132a}, S. Passaggio^{50a}, F. Pastore^{134a,134b,*}, Fr. Pastore⁷⁷, G. Pásztor²⁹, S. Patariaia¹⁷⁵, N.D. Patel¹⁵⁰, J.R. Pater⁸⁴, T. Pauly³⁰, J. Pearce¹⁶⁹, B. Pearson¹¹³, L.E. Pedersen³⁶, M. Pedersen¹¹⁹, S. Pedraza Lopez¹⁶⁷, R. Pedro^{126a,126b}, S.V. Peleganchuk^{109,c}, D. Pelikan¹⁶⁶, H. Peng^{33b}, B. Penning³¹, J. Penwell⁶¹, D.V. Perepelitsa²⁵, E. Perez Codina^{159a}, M.T. Pérez García-Estañ¹⁶⁷, L. Perini^{91a,91b}, H. Pernegger³⁰, S. Perrella^{104a,104b}, R. Peschke⁴², V.D. Peshekhonov⁶⁵, K. Peters³⁰, R.F.Y. Peters⁸⁴, B.A. Petersen³⁰, T.C. Petersen³⁶, E. Petit⁴², A. Petridis^{146a,146b}, C. Petridou¹⁵⁴, E. Petrolo^{132a}, F. Petrucci^{134a,134b}, N.E. Pettersson¹⁵⁷,

R. Pezoa^{32b}, P.W. Phillips¹³¹, G. Piacquadio¹⁴³, E. Pianori¹⁷⁰, A. Picazio⁴⁹, E. Piccaro⁷⁶, M. Piccinini^{20a,20b}, M.A. Pickering¹²⁰, R. Piegai²⁷, D.T. Pignotti¹¹¹, J.E. Pilcher³¹, A.D. Pilkington⁸⁴, J. Pina^{126a,126b,126d}, M. Pinamonti^{164a,164c,ad}, J.L. Pinfold³, A. Pingel³⁶, B. Pinto^{126a}, S. Pires⁸⁰, M. Pitt¹⁷², C. Pizio^{91a,91b}, L. Plazak^{144a}, M.-A. Pleier²⁵, V. Pleskot¹²⁹, E. Plotnikova⁶⁵, P. Plucinski^{146a,146b}, D. Pluth⁶⁴, R. Poettgen⁸³, L. Poggioli¹¹⁷, D. Pohl²¹, G. Polesello^{121a}, A. Policicchio^{37a,37b}, R. Polifka¹⁵⁸, A. Polini^{20a}, C.S. Pollard⁵³, V. Polychronakos²⁵, K. Pommès³⁰, L. Pontecorvo^{132a}, B.G. Pope⁹⁰, G.A. Popeneciu^{26b}, D.S. Popovic¹³, A. Poppleton³⁰, S. Pospisil¹²⁸, K. Potamianos¹⁵, I.N. Potrap⁶⁵, C.J. Potter¹⁴⁹, C.T. Potter¹¹⁶, G. Poulard³⁰, J. Poveda³⁰, V. Pozdnyakov⁶⁵, P. Pralavorio⁸⁵, A. Pranko¹⁵, S. Prasad³⁰, S. Prell⁶⁴, D. Price⁸⁴, L.E. Price⁶, M. Primavera^{73a}, S. Prince⁸⁷, M. Proissl⁴⁶, K. Prokofiev^{60c}, F. Prokoshin^{32b}, E. Protopapadaki¹³⁶, S. Protopopescu²⁵, J. Proudfoot⁶, M. Przybycien^{38a}, E. Ptacek¹¹⁶, D. Puddu^{134a,134b}, E. Pueschel⁸⁶, D. Puldon¹⁴⁸, M. Purohit^{25,ae}, P. Puzo¹¹⁷, J. Qian⁸⁹, G. Qin⁵³, Y. Qin⁸⁴, A. Quadt⁵⁴, D.R. Quarrie¹⁵, W.B. Quayle^{164a,164b}, M. Queitsch-Maitland⁸⁴, D. Quilty⁵³, S. Raddum¹¹⁹, V. Radeka²⁵, V. Radescu⁴², S.K. Radhakrishnan¹⁴⁸, P. Radloff¹¹⁶, P. Rados⁸⁸, F. Ragusa^{91a,91b}, G. Rahal¹⁷⁸, S. Rajagopalan²⁵, M. Rammensee³⁰, C. Rangel-Smith¹⁶⁶, F. Rauscher¹⁰⁰, S. Rave⁸³, T. Ravenscroft⁵³, M. Raymond³⁰, A.L. Read¹¹⁹, N.P. Readioff⁷⁴, D.M. Rebuzzi^{121a,121b}, A. Redelbach¹⁷⁴, G. Redlinger²⁵, R. Reece¹³⁷, K. Reeves⁴¹, L. Rehnisch¹⁶, H. Reisin²⁷, M. Relich¹⁶³, C. Rembser³⁰, H. Ren^{33a}, A. Renaud¹¹⁷, M. Rescigno^{132a}, S. Resconi^{91a}, O.L. Rezanova^{109,c}, P. Reznicek¹²⁹, R. Rezvani⁹⁵, R. Richter¹⁰¹, S. Richter⁷⁸, E. Richter-Was^{38b}, O. Ricken²¹, M. Ridel⁸⁰, P. Rieck¹⁶, C.J. Riegel¹⁷⁵, J. Rieger⁵⁴, M. Rijssenbeek¹⁴⁸, A. Rimoldi^{121a,121b}, L. Rinaldi^{20a}, B. Ristić⁴⁹, E. Ritsch⁶², I. Riu¹², F. Rizatdinova¹¹⁴, E. Rizvi⁷⁶, S.H. Robertson^{87,k}, A. Robichaud-Veronneau⁸⁷, D. Robinson²⁸, J.E.M. Robinson⁸⁴, A. Robson⁵³, C. Roda^{124a,124b}, S. Roe³⁰, O. Røhne¹¹⁹, S. Rolli¹⁶¹, A. Romaniouk⁹⁸, M. Romano^{20a,20b}, S.M. Romano Saez³⁴, E. Romero Adam¹⁶⁷, N. Rompotis¹³⁸, M. Ronzani⁴⁸, L. Roos⁸⁰, E. Ros¹⁶⁷, S. Rosati^{132a}, K. Rosbach⁴⁸, P. Rose¹³⁷, P.L. Rosendahl¹⁴, O. Rosenthal¹⁴¹, V. Rossetti^{146a,146b}, E. Rossi^{104a,104b}, L.P. Rossi^{50a}, R. Rosten¹³⁸, M. Rotaru^{26a}, I. Roth¹⁷², J. Rothberg¹³⁸, D. Rousseau¹¹⁷, C.R. Royon¹³⁶, A. Rozanov⁸⁵, Y. Rozen¹⁵², X. Ruan^{145c}, F. Rubbo¹⁴³, I. Rubinskiy⁴², V.I. Rud⁹⁹, C. Rudolph⁴⁴, M.S. Rudolph¹⁵⁸, F. Rühr⁴⁸, A. Ruiz-Martinez³⁰, Z. Rurikova⁴⁸, N.A. Rusakovich⁶⁵, A. Ruschke¹⁰⁰, H.L. Russell¹³⁸, J.P. Rutherford⁷, N. Ruthmann⁴⁸, Y.F. Ryabov¹²³, M. Rybar¹⁶⁵, G. Rybkin¹¹⁷, N.C. Ryder¹²⁰, A.F. Saavedra¹⁵⁰, G. Sabato¹⁰⁷, S. Sacerdoti²⁷, A. Saddique³, H.F.-W. Sadrozinski¹³⁷, R. Sadykov⁶⁵, F. Safai Tehrani^{132a}, M. Saimpert¹³⁶, H. Sakamoto¹⁵⁵, Y. Sakurai¹⁷¹, G. Salamanna^{134a,134b}, A. Salamon^{133a}, M. Saleem¹¹³, D. Salek¹⁰⁷, P.H. Sales De Bruin¹³⁸, D. Salihagic¹⁰¹, A. Salnikov¹⁴³, J. Salt¹⁶⁷, D. Salvatore^{37a,37b}, F. Salvatore¹⁴⁹, A. Salvucci¹⁰⁶, A. Salzburger³⁰, D. Sampsonidis¹⁵⁴, A. Sanchez^{104a,104b}, J. Sánchez¹⁶⁷, V. Sanchez Martinez¹⁶⁷, H. Sandaker¹¹⁹, R.L. Sandbach⁷⁶, H.G. Sander⁸³, M.P. Sanders¹⁰⁰, M. Sandhoff¹⁷⁵, C. Sandoval¹⁶², R. Sandstroem¹⁰¹, D.P.C. Sankey¹³¹, M. Sannino^{50a,50b}, A. Sansoni⁴⁷, C. Santoni³⁴, R. Santonico^{133a,133b}, H. Santos^{126a}, I. Santoyo Castillo¹⁴⁹, K. Sapp¹²⁵, A. Saprnov⁶⁵, J.G. Saraiva^{126a,126d}, B. Sarrazin²¹, O. Sasaki⁶⁶, Y. Sasaki¹⁵⁵, K. Sato¹⁶⁰, G. Sauvage^{5,*}, E. Sauvan⁵, G. Savage⁷⁷, P. Savard^{158,d}, C. Sawyer¹²⁰, L. Sawyer^{79,n}, J. Saxon³¹, C. Sbarra^{20a}, A. Sbrizzi^{20a,20b}, T. Scanlon⁷⁸, D.A. Scannicchio¹⁶³, M. Scarcella¹⁵⁰, V. Scarfone^{37a,37b}, J. Schaarschmidt¹⁷², P. Schacht¹⁰¹, D. Schaefer³⁰, R. Schaefer⁴², J. Schaeffer⁸³, S. Schaepe²¹, S. Schaetzel^{58b}, U. Schäfer⁸³, A.C. Schaffer¹¹⁷, D. Schaile¹⁰⁰, R.D. Schamberger¹⁴⁸, V. Scharf^{58a}, V.A. Schegelsky¹²³, D. Scheirich¹²⁹, M. Schernau¹⁶³, C. Schiavi^{50a,50b}, C. Schillo⁴⁸, M. Schioppa^{37a,37b}, S. Schlenker³⁰, E. Schmidt⁴⁸, K. Schmieden³⁰, C. Schmitt⁸³, S. Schmitt^{58b}, S. Schmitt⁴², B. Schneider^{159a}, Y.J. Schnellbach⁷⁴, U. Schnoor⁴⁴, L. Schoeffel¹³⁶, A. Schoening^{58b}, B.D. Schoenrock⁹⁰, E. Schopf²¹, A.L.S. Schorlemmer⁵⁴, M. Schott⁸³, D. Schouten^{159a},

J. Schovancova⁸, S. Schramm¹⁵⁸, M. Schreyer¹⁷⁴, C. Schroeder⁸³, N. Schuh⁸³, M.J. Schultens²¹, H.-C. Schultz-Coulon^{58a}, H. Schulz¹⁶, M. Schumacher⁴⁸, B.A. Schumm¹³⁷, Ph. Schune¹³⁶, C. Schwanenberger⁸⁴, A. Schwartzman¹⁴³, T.A. Schwarz⁸⁹, Ph. Schwegler¹⁰¹, H. Schweiger⁸⁴, Ph. Schwemling¹³⁶, R. Schwienhorst⁹⁰, J. Schwindling¹³⁶, T. Schwindt²¹, M. Schwoerer⁵, F.G. Sciacca¹⁷, E. Scifo¹¹⁷, G. Sciolla²³, F. Scuri^{124a,124b}, F. Scutti²¹, J. Searcy⁸⁹, G. Sedov⁴², E. Sedykh¹²³, P. Seema²¹, S.C. Seidel¹⁰⁵, A. Seiden¹³⁷, F. Seifert¹²⁸, J.M. Seixas^{24a}, G. Sekhniaidze^{104a}, K. Sekhon⁸⁹, S.J. Sekula⁴⁰, K.E. Selbach⁴⁶, D.M. Seliverstov^{123,*}, N. Semprini-Cesari^{20a,20b}, C. Serfon³⁰, L. Serin¹¹⁷, L. Serkin^{164a,164b}, T. Serre⁸⁵, M. Sessa^{134a,134b}, R. Seuster^{159a}, H. Severini¹¹³, T. Sfiligoj⁷⁵, F. Sforza¹⁰¹, A. Sfyrlla³⁰, E. Shabalina⁵⁴, M. Shamim¹¹⁶, L.Y. Shan^{33a}, R. Shang¹⁶⁵, J.T. Shank²², M. Shapiro¹⁵, P.B. Shatalov⁹⁷, K. Shaw^{164a,164b}, S.M. Shaw⁸⁴, A. Shcherbakova^{146a,146b}, C.Y. Shehu¹⁴⁹, P. Sherwood⁷⁸, L. Shi^{151,af}, S. Shimizu⁶⁷, C.O. Shimmin¹⁶³, M. Shimojima¹⁰², M. Shiyakova⁶⁵, A. Shmeleva⁹⁶, D. Shoaleh Saadi⁹⁵, M.J. Shochet³¹, S. Shojaii^{91a,91b}, S. Shrestha¹¹¹, E. Shulga⁹⁸, M.A. Shupe⁷, S. Shushkevich⁴², P. Sicho¹²⁷, O. Sidiropoulou¹⁷⁴, D. Sidorov¹¹⁴, A. Sidoti^{20a,20b}, F. Siegert⁴⁴, Dj. Sijacki¹³, J. Silva^{126a,126d}, Y. Silver¹⁵³, S.B. Silverstein^{146a}, V. Simak¹²⁸, O. Simard⁵, Lj. Simic¹³, S. Simion¹¹⁷, E. Simioni⁸³, B. Simmons⁷⁸, D. Simon³⁴, R. Simoniello^{91a,91b}, P. Sinervo¹⁵⁸, N.B. Sinev¹¹⁶, G. Siragusa¹⁷⁴, A.N. Sisakyan^{65,*}, S.Yu. Sivoklokov⁹⁹, J. Sjölin^{146a,146b}, T.B. Sjursen¹⁴, M.B. Skinner⁷², H.P. Skottowe⁵⁷, P. Skubic¹¹³, M. Slater¹⁸, T. Slavicek¹²⁸, M. Slawinska¹⁰⁷, K. Sliwa¹⁶¹, V. Smakhtin¹⁷², B.H. Smart⁴⁶, L. Smestad¹⁴, S.Yu. Smirnov⁹⁸, Y. Smirnov⁹⁸, L.N. Smirnova^{99,ag}, O. Smirnova⁸¹, M.N.K. Smith³⁵, R.W. Smith³⁵, M. Smizanska⁷², K. Smolek¹²⁸, A.A. Snesev⁹⁶, G. Snidero⁷⁶, S. Snyder²⁵, R. Sobie^{169,k}, F. Socher⁴⁴, A. Soffer¹⁵³, D.A. Soh^{151,af}, C.A. Solans³⁰, M. Solar¹²⁸, J. Solc¹²⁸, E.Yu. Soldatov⁹⁸, U. Soldevila¹⁶⁷, A.A. Solodkov¹³⁰, A. Soloshenko⁶⁵, O.V. Solovyanov¹³⁰, V. Solovyev¹²³, P. Sommer⁴⁸, H.Y. Song^{33b}, N. Soni¹, A. Sood¹⁵, A. Sopczak¹²⁸, B. Sopko¹²⁸, V. Sopko¹²⁸, V. Sorin¹², D. Sosa^{58b}, M. Sosebee⁸, C.L. Sotiropoulou^{124a,124b}, R. Soualah^{164a,164c}, P. Soueid⁹⁵, A.M. Soukharev^{109,c}, D. South⁴², B.C. Sowden⁷⁷, S. Spagnolo^{73a,73b}, M. Spalla^{124a,124b}, F. Spanò⁷⁷, W.R. Spearman⁵⁷, F. Spettel¹⁰¹, R. Spighi^{20a}, G. Spigo³⁰, L.A. Spiller⁸⁸, M. Spousta¹²⁹, T. Spreitzer¹⁵⁸, R.D. St. Denis^{53,*}, S. Staerz⁴⁴, J. Stahlman¹²², R. Stamen^{58a}, S. Stamm¹⁶, E. Stanecka³⁹, C. Stanescu^{134a}, M. Stanescu-Bellu⁴², M.M. Stanitzki⁴², S. Stapnes¹¹⁹, E.A. Starchenko¹³⁰, J. Stark⁵⁵, P. Staroba¹²⁷, P. Starovoitov⁴², R. Staszewski³⁹, P. Stavina^{144a,*}, P. Steinberg²⁵, B. Stelzer¹⁴², H.J. Stelzer³⁰, O. Stelzer-Chilton^{159a}, H. Stenzel⁵², S. Stern¹⁰¹, G.A. Stewart⁵³, J.A. Stillings²¹, M.C. Stockton⁸⁷, M. Stoebe⁸⁷, G. Stoicea^{26a}, P. Stolte⁵⁴, S. Stonjek¹⁰¹, A.R. Stradling⁸, A. Straessner⁴⁴, M.E. Stramaglia¹⁷, J. Strandberg¹⁴⁷, S. Strandberg^{146a,146b}, A. Strandlie¹¹⁹, E. Strauss¹⁴³, M. Strauss¹¹³, P. Strizenec^{144b}, R. Ströhmer¹⁷⁴, D.M. Strom¹¹⁶, R. Stroynowski⁴⁰, A. Strubig¹⁰⁶, S.A. Stucci¹⁷, B. Stugu¹⁴, N.A. Styles⁴², D. Su¹⁴³, J. Su¹²⁵, R. Subramaniam⁷⁹, A. Succurro¹², Y. Sugaya¹¹⁸, C. Suhr¹⁰⁸, M. Suk¹²⁸, V.V. Sulin⁹⁶, S. Sultansoy^{4c}, T. Sumida⁶⁸, S. Sun⁵⁷, X. Sun^{33a}, J.E. Sundermann⁴⁸, K. Suruliz¹⁴⁹, G. Susinno^{37a,37b}, M.R. Sutton¹⁴⁹, S. Suzuki⁶⁶, Y. Suzuki⁶⁶, M. Svatos¹²⁷, S. Swedish¹⁶⁸, M. Swiatlowski¹⁴³, I. Sykora^{144a}, T. Sykora¹²⁹, D. Ta⁹⁰, C. Taccini^{134a,134b}, K. Tackmann⁴², J. Taenzer¹⁵⁸, A. Taffard¹⁶³, R. Tahirout^{159a}, N. Taiblum¹⁵³, H. Takai²⁵, R. Takashima⁶⁹, H. Takeda⁶⁷, T. Takeshita¹⁴⁰, Y. Takubo⁶⁶, M. Talby⁸⁵, A.A. Talyshchev^{109,c}, J.Y.C. Tam¹⁷⁴, K.G. Tan⁸⁸, J. Tanaka¹⁵⁵, R. Tanaka¹¹⁷, S. Tanaka⁶⁶, B.B. Tannenwald¹¹¹, N. Tannoury²¹, S. Tapprogge⁸³, S. Tarem¹⁵², F. Tarrade²⁹, G.F. Tartarelli^{91a}, P. Tas¹²⁹, M. Tasevsky¹²⁷, T. Tashiro⁶⁸, E. Tassi^{37a,37b}, A. Tavares Delgado^{126a,126b}, Y. Tayalati^{135d}, F.E. Taylor⁹⁴, G.N. Taylor⁸⁸, W. Taylor^{159b}, F.A. Teischinger³⁰, M. Teixeira Dias Castanheira⁷⁶, P. Teixeira-Dias⁷⁷, K.K. Temming⁴⁸, H. Ten Kate³⁰, P.K. Teng¹⁵¹, J.J. Teoh¹¹⁸, F. Tepel¹⁷⁵, S. Terada⁶⁶, K. Terashi¹⁵⁵, J. Terron⁸², S. Terzo¹⁰¹, M. Testa⁴⁷, R.J. Teuscher^{158,k},

J. Therhaag²¹, T. Theveneaux-Pelzer³⁴, J.P. Thomas¹⁸, J. Thomas-Wilsker⁷⁷, E.N. Thompson³⁵, P.D. Thompson¹⁸, R.J. Thompson⁸⁴, A.S. Thompson⁵³, L.A. Thomsen¹⁷⁶, E. Thomson¹²², M. Thomson²⁸, R.P. Thun^{89,*}, M.J. Tibbetts¹⁵, R.E. Ticse Torres⁸⁵, V.O. Tikhomirov^{96,ah}, Yu.A. Tikhonov^{109,c}, S. Timoshenko⁹⁸, E. Tiouchichine⁸⁵, P. Tipton¹⁷⁶, S. Tisserant⁸⁵, T. Todorov^{5,*}, S. Todorova-Nova¹²⁹, J. Tojo⁷⁰, S. Tokár^{144a}, K. Tokushuku⁶⁶, K. Tollefson⁹⁰, E. Tolley⁵⁷, L. Tomlinson⁸⁴, M. Tomoto¹⁰³, L. Tompkins^{143,ai}, K. Toms¹⁰⁵, E. Torrence¹¹⁶, H. Torres¹⁴², E. Torró Pastor¹⁶⁷, J. Toth^{85,aj}, F. Touchard⁸⁵, D.R. Tovey¹³⁹, T. Trefzger¹⁷⁴, L. Tremblet³⁰, A. Tricoli³⁰, I.M. Trigger^{159a}, S. Trincas-Duvoid⁸⁰, M.F. Tripiana¹², W. Trischuk¹⁵⁸, B. Trocmé⁵⁵, C. Troncon^{91a}, M. Trottier-McDonald¹⁵, M. Trovatelli^{134a,134b}, P. True⁹⁰, L. Truong^{164a,164c}, M. Trzebinski³⁹, A. Trzupek³⁹, C. Tsarouchas³⁰, J.C.-L. Tseng¹²⁰, P.V. Tsiarehsha⁹², D. Tsionou¹⁵⁴, G. Tsipolitis¹⁰, N. Tsirintanis⁹, S. Tsiskaridze¹², V. Tsiskaridze⁴⁸, E.G. Tskhadadze^{51a}, I.I. Tsukerman⁹⁷, V. Tsulaia¹⁵, S. Tsuno⁶⁶, D. Tsybychev¹⁴⁸, A. Tudorache^{26a}, V. Tudorache^{26a}, A.N. Tuna¹²², S.A. Tupputi^{20a,20b}, S. Turchikhin^{99,ag}, D. Turecek¹²⁸, R. Turra^{91a,91b}, A.J. Turvey⁴⁰, P.M. Tuts³⁵, A. Tykhonov⁴⁹, M. Tylmad^{146a,146b}, M. Tyndel¹³¹, I. Ueda¹⁵⁵, R. Ueno²⁹, M. Ughetto^{146a,146b}, M. Ugland¹⁴, M. Uhlenbrock²¹, F. Ukegawa¹⁶⁰, G. Unal³⁰, A. Undrus²⁵, G. Unel¹⁶³, F.C. Ungaro⁴⁸, Y. Unno⁶⁶, C. Unverdorben¹⁰⁰, J. Urban^{144b}, P. Urquijo⁸⁸, P. Urrejola⁸³, G. Usai⁸, A. Usanova⁶², L. Vacavant⁸⁵, V. Vacek¹²⁸, B. Vachon⁸⁷, C. Valderanis⁸³, N. Valencic¹⁰⁷, S. Valentinetti^{20a,20b}, A. Valero¹⁶⁷, L. Valery¹², S. Valkar¹²⁹, E. Valladolid Gallego¹⁶⁷, S. Vallecorsa⁴⁹, J.A. Valls Ferrer¹⁶⁷, W. Van Den Wollenberg¹⁰⁷, P.C. Van Der Deijl¹⁰⁷, R. van der Geer¹⁰⁷, H. van der Graaf¹⁰⁷, R. Van Der Leeuw¹⁰⁷, N. van Eldik¹⁵², P. van Gemmeren⁶, J. Van Nieuwkoop¹⁴², I. van Vulpen¹⁰⁷, M.C. van Woerden³⁰, M. Vanadia^{132a,132b}, W. Vandelli³⁰, R. Vanguri¹²², A. Vaniachine⁶, F. Vannucci⁸⁰, G. Vardanyan¹⁷⁷, R. Vari^{132a}, E.W. Varnes⁷, T. Varol⁴⁰, D. Varouchas⁸⁰, A. Vartapetian⁸, K.E. Varvell¹⁵⁰, F. Vazeille³⁴, T. Vazquez Schroeder⁸⁷, J. Veatch⁷, L.M. Veloce¹⁵⁸, F. Veloso^{126a,126c}, T. Velz²¹, S. Veneziano^{132a}, A. Ventura^{73a,73b}, D. Ventura⁸⁶, M. Venturi¹⁶⁹, N. Venturi¹⁵⁸, A. Venturini²³, V. Vercesi^{121a}, M. Verducci^{132a,132b}, W. Verkerke¹⁰⁷, J.C. Vermeulen¹⁰⁷, A. Vest⁴⁴, M.C. Vetterli^{142,d}, O. Viazlo⁸¹, I. Vichou¹⁶⁵, T. Vickey¹³⁹, O.E. Vickey Boeriu¹³⁹, G.H.A. Viehhauser¹²⁰, S. Viel¹⁵, R. Vigne³⁰, M. Villa^{20a,20b}, M. Villaplana Perez^{91a,91b}, E. Vilucchi⁴⁷, M.G. Vincker²⁹, V.B. Vinogradov⁶⁵, I. Vivarelli¹⁴⁹, F. Vives Vaque³, S. Vlachos¹⁰, D. Vladoiu¹⁰⁰, M. Vlasak¹²⁸, M. Vogel^{32a}, P. Vokac¹²⁸, G. Volpi^{124a,124b}, M. Volpi⁸⁸, H. von der Schmitt¹⁰¹, H. von Radziewski⁴⁸, E. von Toerne²¹, V. Vorobel¹²⁹, K. Vorobev⁹⁸, M. Vos¹⁶⁷, R. Voss³⁰, J.H. Vosseveld⁷⁴, N. Vranjes¹³, M. Vranjes Milosavljevic¹³, V. Vrba¹²⁷, M. Vreeswijk¹⁰⁷, R. Vuillermet³⁰, I. Vukotic³¹, Z. Vykydal¹²⁸, P. Wagner²¹, W. Wagner¹⁷⁵, H. Wahlberg⁷¹, S. Wahrmund⁴⁴, J. Wakabayashi¹⁰³, J. Walder⁷², R. Walker¹⁰⁰, W. Walkowiak¹⁴¹, C. Wang^{33c}, F. Wang¹⁷³, H. Wang¹⁵, H. Wang⁴⁰, J. Wang⁴², J. Wang^{33a}, K. Wang⁸⁷, R. Wang⁶, S.M. Wang¹⁵¹, T. Wang²¹, X. Wang¹⁷⁶, C. Wanotayaroj¹¹⁶, A. Warburton⁸⁷, C.P. Ward²⁸, D.R. Wardrope⁷⁸, M. Warsinsky⁴⁸, A. Washbrook⁴⁶, C. Wasicki⁴², P.M. Watkins¹⁸, A.T. Watson¹⁸, I.J. Watson¹⁵⁰, M.F. Watson¹⁸, G. Watts¹³⁸, S. Watts⁸⁴, B.M. Waugh⁷⁸, S. Webb⁸⁴, M.S. Weber¹⁷, S.W. Weber¹⁷⁴, J.S. Webster³¹, A.R. Weidberg¹²⁰, B. Weinert⁶¹, J. Weingarten⁵⁴, C. Weiser⁴⁸, H. Weits¹⁰⁷, P.S. Wells³⁰, T. Wenaus²⁵, T. Wengler³⁰, S. Wenig³⁰, N. Wermes²¹, M. Werner⁴⁸, P. Werner³⁰, M. Wessels^{58a}, J. Wetter¹⁶¹, K. Whalen²⁹, A.M. Wharton⁷², A. White⁸, M.J. White¹, R. White^{32b}, S. White^{124a,124b}, D. Whiteson¹⁶³, F.J. Wickens¹³¹, W. Wiedenmann¹⁷³, M. Wielers¹³¹, P. Wienemann²¹, C. Wiglesworth³⁶, L.A.M. Wiik-Fuchs²¹, A. Wildauer¹⁰¹, H.G. Wilkens³⁰, H.H. Williams¹²², S. Williams¹⁰⁷, C. Willis⁹⁰, S. Willocq⁸⁶, A. Wilson⁸⁹, J.A. Wilson¹⁸, I. Wingerter-Seez⁵, F. Winklmeier¹¹⁶, B.T. Winter²¹, M. Wittgen¹⁴³, J. Wittkowski¹⁰⁰, S.J. Wollstadt⁸³, M.W. Wolter³⁹, H. Wolters^{126a,126c}, B.K. Wosiek³⁹, J. Wotschack³⁰, M.J. Woudstra⁸⁴, K.W. Wozniak³⁹, M. Wu⁵⁵,

M. Wu³¹, S.L. Wu¹⁷³, X. Wu⁴⁹, Y. Wu⁸⁹, T.R. Wyatt⁸⁴, B.M. Wynne⁴⁶, S. Xella³⁶, D. Xu^{33a}, L. Xu^{33b,ak}, B. Yabsley¹⁵⁰, S. Yacoob^{145b,al}, R. Yakabe⁶⁷, M. Yamada⁶⁶, Y. Yamaguchi¹¹⁸, A. Yamamoto⁶⁶, S. Yamamoto¹⁵⁵, T. Yamanaka¹⁵⁵, K. Yamauchi¹⁰³, Y. Yamazaki⁶⁷, Z. Yan²², H. Yang^{33e}, H. Yang¹⁷³, Y. Yang¹⁵¹, L. Yao^{33a}, W.-M. Yao¹⁵, Y. Yasu⁶⁶, E. Yatsenko⁵, K.H. Yau Wong²¹, J. Ye⁴⁰, S. Ye²⁵, I. Yeletsikh⁶⁵, A.L. Yen⁵⁷, E. Yildirim⁴², K. Yorita¹⁷¹, R. Yoshida⁶, K. Yoshihara¹²², C. Young¹⁴³, C.J.S. Young³⁰, S. Youssef²², D.R. Yu¹⁵, J. Yu⁸, J.M. Yu⁸⁹, J. Yu¹¹⁴, L. Yuan⁶⁷, A. Yurkewicz¹⁰⁸, I. Yusuf^{28,am}, B. Zabinski³⁹, R. Zaidan⁶³, A.M. Zaitsev^{130,ab}, J. Zalieckas¹⁴, A. Zaman¹⁴⁸, S. Zambito⁵⁷, L. Zanello^{132a,132b}, D. Zanzi⁸⁸, C. Zeitnitz¹⁷⁵, M. Zeman¹²⁸, A. Zemla^{38a}, K. Zengel²³, O. Zenin¹³⁰, T. Ženiš^{144a}, D. Zerwas¹¹⁷, D. Zhang⁸⁹, F. Zhang¹⁷³, J. Zhang⁶, L. Zhang⁴⁸, R. Zhang^{33b}, X. Zhang^{33d}, Z. Zhang¹¹⁷, X. Zhao⁴⁰, Y. Zhao^{33d,117}, Z. Zhao^{33b}, A. Zhemchugov⁶⁵, J. Zhong¹²⁰, B. Zhou⁸⁹, C. Zhou⁴⁵, L. Zhou³⁵, L. Zhou⁴⁰, N. Zhou¹⁶³, C.G. Zhu^{33d}, H. Zhu^{33a}, J. Zhu⁸⁹, Y. Zhu^{33b}, X. Zhuang^{33a}, K. Zhukov⁹⁶, A. Zibell¹⁷⁴, D. Zieminska⁶¹, N.I. Zimine⁶⁵, C. Zimmermann⁸³, S. Zimmermann⁴⁸, Z. Zinonos⁵⁴, M. Zinser⁸³, M. Ziolkowski¹⁴¹, L. Živković¹³, G. Zobernig¹⁷³, A. Zoccoli^{20a,20b}, M. zur Nedden¹⁶, G. Zurzolo^{104a,104b}, L. Zwalinski³⁰.

¹ Department of Physics, University of Adelaide, Adelaide, Australia

² Physics Department, SUNY Albany, Albany NY, United States of America

³ Department of Physics, University of Alberta, Edmonton AB, Canada

⁴ ^(a) Department of Physics, Ankara University, Ankara; ^(b) Istanbul Aydin University, Istanbul; ^(c) Division of Physics, TOBB University of Economics and Technology, Ankara, Turkey

⁵ LAPP, CNRS/IN2P3 and Université Savoie Mont Blanc, Annecy-le-Vieux, France

⁶ High Energy Physics Division, Argonne National Laboratory, Argonne IL, United States of America

⁷ Department of Physics, University of Arizona, Tucson AZ, United States of America

⁸ Department of Physics, The University of Texas at Arlington, Arlington TX, United States of America

⁹ Physics Department, University of Athens, Athens, Greece

¹⁰ Physics Department, National Technical University of Athens, Zografou, Greece

¹¹ Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan

¹² Institut de Física d'Altes Energies and Departament de Física de la Universitat Autònoma de Barcelona, Barcelona, Spain

¹³ Institute of Physics, University of Belgrade, Belgrade, Serbia

¹⁴ Department for Physics and Technology, University of Bergen, Bergen, Norway

¹⁵ Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley CA, United States of America

¹⁶ Department of Physics, Humboldt University, Berlin, Germany

¹⁷ Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland

¹⁸ School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom

¹⁹ ^(a) Department of Physics, Bogazici University, Istanbul; ^(b) Department of Physics Engineering, Gaziantep University, Gaziantep; ^(c) Department of Physics, Dogus University, Istanbul, Turkey

²⁰ ^(a) INFN Sezione di Bologna; ^(b) Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy

²¹ Physikalisches Institut, University of Bonn, Bonn, Germany

²² Department of Physics, Boston University, Boston MA, United States of America

²³ Department of Physics, Brandeis University, Waltham MA, United States of America

²⁴ ^(a) Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro; ^(b) Electrical Circuits Department, Federal University of Juiz de Fora (UFJF), Juiz de Fora; ^(c) Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei; ^(d) Instituto de Física, Universidade de Sao Paulo, Sao Paulo, Brazil

²⁵ Physics Department, Brookhaven National Laboratory, Upton NY, United States of America

- ²⁶ ^(a) *National Institute of Physics and Nuclear Engineering, Bucharest;* ^(b) *National Institute for Research and Development of Isotopic and Molecular Technologies, Physics Department, Cluj Napoca;* ^(c) *University Politehnica Bucharest, Bucharest;* ^(d) *West University in Timisoara, Timisoara, Romania*
- ²⁷ *Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina*
- ²⁸ *Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom*
- ²⁹ *Department of Physics, Carleton University, Ottawa ON, Canada*
- ³⁰ *CERN, Geneva, Switzerland*
- ³¹ *Enrico Fermi Institute, University of Chicago, Chicago IL, United States of America*
- ³² ^(a) *Departamento de Física, Pontificia Universidad Católica de Chile, Santiago;* ^(b) *Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile*
- ³³ ^(a) *Institute of High Energy Physics, Chinese Academy of Sciences, Beijing;* ^(b) *Department of Modern Physics, University of Science and Technology of China, Anhui;* ^(c) *Department of Physics, Nanjing University, Jiangsu;* ^(d) *School of Physics, Shandong University, Shandong;* ^(e) *Department of Physics and Astronomy, Shanghai Key Laboratory for Particle Physics and Cosmology, Shanghai Jiao Tong University, Shanghai;* ^(f) *Physics Department, Tsinghua University, Beijing 100084, China*
- ³⁴ *Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France*
- ³⁵ *Nevis Laboratory, Columbia University, Irvington NY, United States of America*
- ³⁶ *Niels Bohr Institute, University of Copenhagen, Kobenhavn, Denmark*
- ³⁷ ^(a) *INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati;* ^(b) *Dipartimento di Fisica, Università della Calabria, Rende, Italy*
- ³⁸ ^(a) *AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow;* ^(b) *Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland*
- ³⁹ *Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland*
- ⁴⁰ *Physics Department, Southern Methodist University, Dallas TX, United States of America*
- ⁴¹ *Physics Department, University of Texas at Dallas, Richardson TX, United States of America*
- ⁴² *DESY, Hamburg and Zeuthen, Germany*
- ⁴³ *Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany*
- ⁴⁴ *Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany*
- ⁴⁵ *Department of Physics, Duke University, Durham NC, United States of America*
- ⁴⁶ *SUPA - School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom*
- ⁴⁷ *INFN Laboratori Nazionali di Frascati, Frascati, Italy*
- ⁴⁸ *Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany*
- ⁴⁹ *Section de Physique, Université de Genève, Geneva, Switzerland*
- ⁵⁰ ^(a) *INFN Sezione di Genova;* ^(b) *Dipartimento di Fisica, Università di Genova, Genova, Italy*
- ⁵¹ ^(a) *E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi;* ^(b) *High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia*
- ⁵² *II Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany*
- ⁵³ *SUPA - School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom*
- ⁵⁴ *II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany*
- ⁵⁵ *Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS/IN2P3, Grenoble, France*
- ⁵⁶ *Department of Physics, Hampton University, Hampton VA, United States of America*
- ⁵⁷ *Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge MA, United States of America*
- ⁵⁸ ^(a) *Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg;* ^(b) *Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg;* ^(c) *ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany*
- ⁵⁹ *Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan*

- ⁶⁰ ^(a) Department of Physics, The Chinese University of Hong Kong, Shatin, N.T., Hong Kong; ^(b) Department of Physics, The University of Hong Kong, Hong Kong; ^(c) Department of Physics, The Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China
- ⁶¹ Department of Physics, Indiana University, Bloomington IN, United States of America
- ⁶² Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
- ⁶³ University of Iowa, Iowa City IA, United States of America
- ⁶⁴ Department of Physics and Astronomy, Iowa State University, Ames IA, United States of America
- ⁶⁵ Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia
- ⁶⁶ KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
- ⁶⁷ Graduate School of Science, Kobe University, Kobe, Japan
- ⁶⁸ Faculty of Science, Kyoto University, Kyoto, Japan
- ⁶⁹ Kyoto University of Education, Kyoto, Japan
- ⁷⁰ Department of Physics, Kyushu University, Fukuoka, Japan
- ⁷¹ Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
- ⁷² Physics Department, Lancaster University, Lancaster, United Kingdom
- ⁷³ ^(a) INFN Sezione di Lecce; ^(b) Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy
- ⁷⁴ Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom
- ⁷⁵ Department of Physics, Jožef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia
- ⁷⁶ School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom
- ⁷⁷ Department of Physics, Royal Holloway University of London, Surrey, United Kingdom
- ⁷⁸ Department of Physics and Astronomy, University College London, London, United Kingdom
- ⁷⁹ Louisiana Tech University, Ruston LA, United States of America
- ⁸⁰ Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France
- ⁸¹ Fysiska institutionen, Lunds universitet, Lund, Sweden
- ⁸² Departamento de Física Teórica C-15, Universidad Autónoma de Madrid, Madrid, Spain
- ⁸³ Institut für Physik, Universität Mainz, Mainz, Germany
- ⁸⁴ School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
- ⁸⁵ CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
- ⁸⁶ Department of Physics, University of Massachusetts, Amherst MA, United States of America
- ⁸⁷ Department of Physics, McGill University, Montreal QC, Canada
- ⁸⁸ School of Physics, University of Melbourne, Victoria, Australia
- ⁸⁹ Department of Physics, The University of Michigan, Ann Arbor MI, United States of America
- ⁹⁰ Department of Physics and Astronomy, Michigan State University, East Lansing MI, United States of America
- ⁹¹ ^(a) INFN Sezione di Milano; ^(b) Dipartimento di Fisica, Università di Milano, Milano, Italy
- ⁹² B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Republic of Belarus
- ⁹³ National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Republic of Belarus
- ⁹⁴ Department of Physics, Massachusetts Institute of Technology, Cambridge MA, United States of America
- ⁹⁵ Group of Particle Physics, University of Montreal, Montreal QC, Canada
- ⁹⁶ P.N. Lebedev Institute of Physics, Academy of Sciences, Moscow, Russia
- ⁹⁷ Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia
- ⁹⁸ National Research Nuclear University MEPhI, Moscow, Russia
- ⁹⁹ D.V. Skobeltsyn Institute of Nuclear Physics, M.V. Lomonosov Moscow State University, Moscow, Russia
- ¹⁰⁰ Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
- ¹⁰¹ Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
- ¹⁰² Nagasaki Institute of Applied Science, Nagasaki, Japan

- 103 Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan
- 104 (a) INFN Sezione di Napoli; (b) Dipartimento di Fisica, Università di Napoli, Napoli, Italy
- 105 Department of Physics and Astronomy, University of New Mexico, Albuquerque NM, United States of America
- 106 Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
- 107 Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
- 108 Department of Physics, Northern Illinois University, DeKalb IL, United States of America
- 109 Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia
- 110 Department of Physics, New York University, New York NY, United States of America
- 111 Ohio State University, Columbus OH, United States of America
- 112 Faculty of Science, Okayama University, Okayama, Japan
- 113 Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman OK, United States of America
- 114 Department of Physics, Oklahoma State University, Stillwater OK, United States of America
- 115 Palacký University, RCPTM, Olomouc, Czech Republic
- 116 Center for High Energy Physics, University of Oregon, Eugene OR, United States of America
- 117 LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France
- 118 Graduate School of Science, Osaka University, Osaka, Japan
- 119 Department of Physics, University of Oslo, Oslo, Norway
- 120 Department of Physics, Oxford University, Oxford, United Kingdom
- 121 (a) INFN Sezione di Pavia; (b) Dipartimento di Fisica, Università di Pavia, Pavia, Italy
- 122 Department of Physics, University of Pennsylvania, Philadelphia PA, United States of America
- 123 National Research Centre “Kurchatov Institute” B.P.Konstantinov Petersburg Nuclear Physics Institute, St. Petersburg, Russia
- 124 (a) INFN Sezione di Pisa; (b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
- 125 Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh PA, United States of America
- 126 (a) Laboratório de Instrumentação e Física Experimental de Partículas - LIP, Lisboa; (b) Faculdade de Ciências, Universidade de Lisboa, Lisboa; (c) Department of Physics, University of Coimbra, Coimbra; (d) Centro de Física Nuclear da Universidade de Lisboa, Lisboa; (e) Departamento de Física, Universidade do Minho, Braga; (f) Departamento de Física Teórica y del Cosmos and CAFPE, Universidad de Granada, Granada (Spain); (g) Dep Física and CEFITEC of Faculdade de Ciências e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal
- 127 Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic
- 128 Czech Technical University in Prague, Praha, Czech Republic
- 129 Faculty of Mathematics and Physics, Charles University in Prague, Praha, Czech Republic
- 130 State Research Center Institute for High Energy Physics, Protvino, Russia
- 131 Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
- 132 (a) INFN Sezione di Roma; (b) Dipartimento di Fisica, Sapienza Università di Roma, Roma, Italy
- 133 (a) INFN Sezione di Roma Tor Vergata; (b) Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
- 134 (a) INFN Sezione di Roma Tre; (b) Dipartimento di Matematica e Fisica, Università Roma Tre, Roma, Italy
- 135 (a) Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies - Université Hassan II, Casablanca; (b) Centre National de l’Energie des Sciences Techniques Nucleaires, Rabat; (c) Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA-Marrakech; (d) Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda; (e) Faculté des sciences, Université Mohammed V-Agdal, Rabat, Morocco
- 136 DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l’Univers), CEA Saclay (Commissariat à l’Energie Atomique et aux Energies Alternatives), Gif-sur-Yvette, France

- 137 *Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz CA, United States of America*
 138 *Department of Physics, University of Washington, Seattle WA, United States of America*
 139 *Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom*
 140 *Department of Physics, Shinshu University, Nagano, Japan*
 141 *Fachbereich Physik, Universität Siegen, Siegen, Germany*
 142 *Department of Physics, Simon Fraser University, Burnaby BC, Canada*
 143 *SLAC National Accelerator Laboratory, Stanford CA, United States of America*
 144 ^(a) *Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava;* ^(b) *Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic*
 145 ^(a) *Department of Physics, University of Cape Town, Cape Town;* ^(b) *Department of Physics, University of Johannesburg, Johannesburg;* ^(c) *School of Physics, University of the Witwatersrand, Johannesburg, South Africa*
 146 ^(a) *Department of Physics, Stockholm University;* ^(b) *The Oskar Klein Centre, Stockholm, Sweden*
 147 *Physics Department, Royal Institute of Technology, Stockholm, Sweden*
 148 *Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook NY, United States of America*
 149 *Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom*
 150 *School of Physics, University of Sydney, Sydney, Australia*
 151 *Institute of Physics, Academia Sinica, Taipei, Taiwan*
 152 *Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel*
 153 *Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel*
 154 *Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece*
 155 *International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan*
 156 *Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan*
 157 *Department of Physics, Tokyo Institute of Technology, Tokyo, Japan*
 158 *Department of Physics, University of Toronto, Toronto ON, Canada*
 159 ^(a) *TRIUMF, Vancouver BC;* ^(b) *Department of Physics and Astronomy, York University, Toronto ON, Canada*
 160 *Faculty of Pure and Applied Sciences, University of Tsukuba, Tsukuba, Japan*
 161 *Department of Physics and Astronomy, Tufts University, Medford MA, United States of America*
 162 *Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia*
 163 *Department of Physics and Astronomy, University of California Irvine, Irvine CA, United States of America*
 164 ^(a) *INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine;* ^(b) *ICTP, Trieste;* ^(c) *Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy*
 165 *Department of Physics, University of Illinois, Urbana IL, United States of America*
 166 *Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden*
 167 *Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain*
 168 *Department of Physics, University of British Columbia, Vancouver BC, Canada*
 169 *Department of Physics and Astronomy, University of Victoria, Victoria BC, Canada*
 170 *Department of Physics, University of Warwick, Coventry, United Kingdom*
 171 *Waseda University, Tokyo, Japan*
 172 *Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel*
 173 *Department of Physics, University of Wisconsin, Madison WI, United States of America*
 174 *Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany*
 175 *Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany*

- ¹⁷⁶ *Department of Physics, Yale University, New Haven CT, United States of America*
- ¹⁷⁷ *Yerevan Physics Institute, Yerevan, Armenia*
- ¹⁷⁸ *Centre de Calcul de l'Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France*
- ^a *Also at Department of Physics, King's College London, London, United Kingdom*
- ^b *Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan*
- ^c *Also at Novosibirsk State University, Novosibirsk, Russia*
- ^d *Also at TRIUMF, Vancouver BC, Canada*
- ^e *Also at Department of Physics, California State University, Fresno CA, United States of America*
- ^f *Also at Department of Physics, University of Fribourg, Fribourg, Switzerland*
- ^g *Also at Departamento de Física e Astronomia, Faculdade de Ciencias, Universidade do Porto, Portugal*
- ^h *Also at Tomsk State University, Tomsk, Russia*
- ⁱ *Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France*
- ^j *Also at Università di Napoli Parthenope, Napoli, Italy*
- ^k *Also at Institute of Particle Physics (IPP), Canada*
- ^l *Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom*
- ^m *Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia*
- ⁿ *Also at Louisiana Tech University, Ruston LA, United States of America*
- ^o *Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain*
- ^p *Also at Department of Physics, National Tsing Hua University, Taiwan*
- ^q *Also at Department of Physics, The University of Texas at Austin, Austin TX, United States of America*
- ^r *Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia*
- ^s *Also at CERN, Geneva, Switzerland*
- ^t *Also at Georgian Technical University (GTU), Tbilisi, Georgia*
- ^u *Also at Ochadai Academic Production, Ochanomizu University, Tokyo, Japan*
- ^v *Also at Manhattan College, New York NY, United States of America*
- ^w *Also at Hellenic Open University, Patras, Greece*
- ^x *Also at Institute of Physics, Academia Sinica, Taipei, Taiwan*
- ^y *Also at LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France*
- ^z *Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan*
- ^{aa} *Also at School of Physics, Shandong University, Shandong, China*
- ^{ab} *Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia*
- ^{ac} *Also at section de Physique, Université de Genève, Geneva, Switzerland*
- ^{ad} *Also at International School for Advanced Studies (SISSA), Trieste, Italy*
- ^{ae} *Also at Department of Physics and Astronomy, University of South Carolina, Columbia SC, United States of America*
- ^{af} *Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China*
- ^{ag} *Also at Faculty of Physics, M.V.Lomonosov Moscow State University, Moscow, Russia*
- ^{ah} *Also at National Research Nuclear University MEPhI, Moscow, Russia*
- ^{ai} *Also at Department of Physics, Stanford University, Stanford CA, United States of America*
- ^{aj} *Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary*
- ^{ak} *Also at Department of Physics, The University of Michigan, Ann Arbor MI, United States of America*
- ^{al} *Also at Discipline of Physics, University of KwaZulu-Natal, Durban, South Africa*
- ^{am} *Also at University of Malaya, Department of Physics, Kuala Lumpur, Malaysia*
- ^{*} *Deceased*